

Optical solitons with Kudryashov's form of local and non-local self-phase modulation structure having fractional temporal evolution

M.A.S. Murad¹, F.K. Hamasalh², A.H. Arnous^{3,4}, A. Biswas^{5,6,7}, A.J.M. Jawad⁸, Y. Yildirim^{9,10}, L. Moraru^{11,12}, C. Dragomir¹³

¹Department of Mathematics, College of Science, University of Duhok, Duhok, Iraq

²Bakrajo Technical Institute, Sulaimani Polytechnic University, Sulaymaniyah, Iraq

³Department of Mathematical Sciences, Saveetha School of Engineering, SIMATS, Chennai 6021, Tamilnadu, India

⁴Research Center of Applied Mathematics, Khazar University, Baku AZ—1096, Azerbaijan

⁵Department of Mathematics and Physics, Grambling State University, Grambling, LA 71245–2715, USA

⁶Department of Physics and Electronics, Khazar University, Baku AZ—1096, Azerbaijan

⁷Department of Mathematics and Applied Mathematics, Sefako Makgatho Health Sciences University, Medunsa–0204, South Africa

⁸Department of Computer Technical Engineering, Al–Rafidain University College, Baghdad–10064, Iraq

⁹Department of Computer Engineering, Biruni University, Istanbul–34010, Turkey

¹⁰Mathematics Research Center, Near East University, 99138 Nicosia, Cyprus

¹¹Faculty of Sciences and Environment, Department of Chemistry, Physics and Environment, Dunarea de Jos University of Galati, 47 Domneasca Street, 800008 Galati, Romania

¹²Department of Physics, Sefako Makgatho Health Sciences University, Medunsa–0204, South Africa

¹³Department of Applied Sciences, Cross–Border Faculty of Humanities, Economics and Engineering, Dunarea de Jos University of Galati, 111 Domneasca Street, Galati–800201, Romania

*Corresponding author e-mail: biswas.anjan@gmail.com

Abstract. The current paper recovers optical soliton solutions with Kudrashov's proposed self-phase modulation structure and with fractional temporal evolution. This model could control and mitigate the Internet bottleneck effect, a severe hindrance to Internet traffic flow across intercontinental distances. Two independent integration approaches have made this retrieval possible. A wide spectrum of soliton solutions emerged with the collective applications of the integration schemes. The parameter constraints for the existence of these solitons are also presented.

Keywords: solitons, Internet bottleneck, integrability.

<https://doi.org/10.15407/spqeo29.01.080>

PACS 02.30.Ik, 42.65.Tg, 42.65.Jx, 89.20.Hh

Manuscript received 15.02.25; revised version received 08.12.25; accepted for publication 18.03.26; published online 25.03.26.

1. Introduction

Optical solitons have made gigantic advances since their first inception in the telecommunications industry more than half a century ago. These solitons serve as bit carriers in optical fibers for transcontinental and transoceanic distances across the globe. Today's total dependence on internet communications across the globe would not have been possible without the unique engineering marvel of soliton science and technology. Therefore, it is imperative to address this technology from all angles to achieve its performance enhancement. One of the problems that has plagued the telecommunication industry is the Internet bottleneck effect [1–5]. This problem needs to be addressed and tackled with

efficiency. Several mathematical measures and means have been proposed to address this unwanted issue [6–10].

One of them is to include spatio-temporal dispersion in addition to the pre-existing chromatic dispersion (CD) [11–15]. This can slow down the solitons across a junction point, and the traffic flow of signals can be monitored uniformly, like street traffic flow. Another measure that has been adopted is the consideration of the time-dependent coefficient of CD and self-phase modulation (SPM) [1]. This can also slow down the soliton, and in its limit, the solitons can be made to halt, allowing the cross-traffic internet to flow smoothly. The traffic light effect can similarly be adopted here as well. The third mechanism is based on

using the fractional temporal evolution, which can slow down the evolution of pulses along one direction, and thus the traffic light effect can be implemented to control the bottleneck effect [16–20]. The current paper addresses this third measure to study the slow evolution of solitons for the governing nonlinear Schrödinger equation (NLSE) with fractional temporal evolution and linear CD. However, the SPM structure stems from the combination of Kudryashov’s quadruple nonlinear effect and dual-form of nonlocal nonlinearity [21–24]. The nonlinearity parameter dictates the arbitrary intensity of pulses. The details of the model and the derivation of the soliton solutions are jotted, pictured, and exhibited in the rest of the paper after a succinct introduction to the model.

1.1. Governing model

This study investigates the time-fractional nonlinear Schrödinger equation, which includes Kudryashov’s arbitrary refractive index alongside two distinct nonlocal nonlinearities [25]:

$$\frac{\mu q}{t^\mu} + aq_{xx} + (c_1|q|^n + c_2|q|^{2n} + c_3|q|^{3n} + c_4|q|^{4n})q + (c_5(|q|^n)_{xx} + c_6(|q|^{2n})_{xx})q = 0, \quad (1)$$

$0 < \mu \leq 1$,

where $q(x, t)$ represents the wave profile and $\frac{\partial^\mu(\cdot)}{\partial t^\mu}$ denotes the conformable derivative, with n being the nonlinearity parameter. The coefficients c_i for $i = 0, 1, 2, 3, 4$ describe the nonlinearity effects, contributing to SPM. Furthermore, the coefficients c_5 and c_6 characterize a generalized non-local formulation of the refractive index. The power-law parameter n represents arbitrary intensity.

In this paper, we aim to implement two different approaches to improve existing techniques for constructing various soliton solutions dealing with NLSE. The modified simplest equation method is a highly effective and straightforward approach for finding solutions to nonlinear partial differential equations. Further, the research shows that this and the Kudryashov method are highly efficient, effective, impactful, straightforward, and easy to apply to various nonlinear partial differential equations.

Definition 1.1. Let $p: (0, \infty) \rightarrow R$. The conformable derivative of order β can be introduced as follows:

$$L_\beta(p)(x) = \lim_{a \rightarrow 0} \frac{p(x+ax^{1-\beta})-p(x)}{a}, \quad (2)$$

for all $x > 0$ and $\beta \in (0, 1]$ [26].

2. Formulation of the problem

Here, we start with inserting the following wave transform into equation (1):

$$f(x, t) = U(\zeta)e^{i\vartheta(x, t)}, \quad (3)$$

$$\zeta = x - v\frac{t^\mu}{\mu}, \vartheta(x, t) = u - kx + w\frac{t^\mu}{\mu},$$

where u represents the phase center, k stands for the soliton frequency, and w denotes the wave frequency. Upon incorporating the transformations from equation (3) into equation (1), the resulting expression is as follows:

$$aU'' + 2ac_6nU^{2n}U'' + c_5nU^nU'' + c_5n(n-1)U^{n-1}U'^2(2bk^2 + w)U + 2c_6n(2n-1)U^{2n-1}U'^2 - c_1U^{n+1} - c_2U^{2n+1} - c_4U^{4n+1} - c_3U^{3n+1} = 0, \quad (4)$$

and the imaginary part

$$(v + 2bk)U' = 0. \quad (5)$$

From the above equation, the velocity has the following form:

$$v = -2ak. \quad (6)$$

Consider

$$U(\zeta) = V^{\frac{1}{n}}(\zeta). \quad (7)$$

Thus, from equations (4) and (7), we obtain

$$a(1-n)V'^2 + nV V'' + 2c_6n^2aV''V^3 + 2c_6n^2aV'^2V^2 + c_5n^2V''V^2 - n^2(2ak^2 + w)V^2 + c_1n^2V^3 + c_2n^2V^4 + c_3n^2V^5 + c_4n^2V^6 = 0. \quad (8)$$

3. Modified simplest equation method

In this section, we present several novel conformable optical soliton solutions for the current model, derived using the modified simplest equation method. We assume that the solution to equation (10) can be expressed as the following series.

$$V(\zeta) = \sum_{i=0}^M f_i G(\zeta)^i, \quad (9)$$

where f_0, f_1, \dots, f_M are unknown constants, and M is a balancing parameter. In equation (8), implementing the balancing principle between V^3V'' and V^6 leads to $M = 1$. Here, from equation (9) the following is obtained:

$$V(\zeta) = f_0 + f_1 G(\zeta), \quad (10)$$

where $G(\zeta)$ satisfies the following equation:

$$G(\zeta)' = h_0 + G(\zeta)^2. \quad (11)$$

Now, the solutions of equation (13) with parameter s can be defined as follows [27]:

If $h_0 < 0$, we get:

$$G_1(\zeta) = -\sqrt{-h_0} \tanh\left(\sqrt{-h_0}(\zeta + s)\right), \quad (12)$$

$$G_2(\zeta) = -\sqrt{-h_0} \coth\left(\sqrt{-h_0}(\zeta + s)\right), \quad (13)$$

$$G_3(\zeta) = \sqrt{-h_0} \left(\begin{array}{l} -\tanh\left(2\sqrt{-h_0}(\zeta + s)\right) \\ \pm \operatorname{sech}\left(2\sqrt{-h_0}(\zeta + s)\right) \end{array} \right), \quad (14)$$

$$G_4(\zeta) = \sqrt{-h_0} \left(\begin{array}{l} -\coth \left(2\sqrt{-h_0}(\zeta + l) \right) \\ \pm \operatorname{csch} \left(2\sqrt{-h_0}(\zeta + s) \right) \end{array} \right), \quad (15)$$

$$G_5(\zeta) = -\frac{\sqrt{-h_0}}{2} \left(\begin{array}{l} \tanh \left(\frac{\sqrt{-h_0}}{2}(\zeta + s) \right) \\ \pm \coth \left(\frac{\sqrt{-h_0}}{2}(\zeta + s) \right) \end{array} \right). \quad (16)$$

Substituting equations (10) and (11) into equation (8) yields a polynomial expressed in terms of powers of $G(\zeta)$. We then arrange terms with similar powers and set each corresponding coefficient to zero. This process generates the following system of algebraic equations:

$$(G(\zeta))^0: -ak^2n^2f_0^2 - n^2wf_0^2 + n^2c_1f_0^3 + n^2c_2f_0^4 + n^2c_3f_0^5 + n^2c_4f_0^6 + af_1^2h_0^2 - an f_1^2 h_0^2 + 2n^2c_6f_0^2f_1^2h_0^2 = 0, \quad (17)$$

$$(G(\zeta))^1: -2ak^2n^2f_0f_1 - 2n^2wf_0f_1 + 3n^2c_1f_0^2f_1 + 4n^2c_2f_0^3f_1 + 5n^2c_3f_0^4f_1 + 6n^2c_4f_0^5f_1 + 2anf_0f_1h_0 + 2n^2c_5f_0^2f_1h_0 + 4n^2c_6f_0^3f_1h_0 + 4n^2c_6f_0f_1^3h_0^2 = 0, \quad (18)$$

$$(G(\zeta))^2: -ak^2n^2f_1^2 - n^2wf_1^2 + 3n^2c_1f_0f_1^2 + 6n^2c_2f_0^2f_1^2 + 10n^2c_3f_0^3f_1^2 + 15n^2c_4f_0^4f_1^2 + 2af_1^2h_0 + 4n^2c_5f_0f_1^2h_0 + 16n^2c_6f_0^2f_1^2h_0 + 2n^2c_6f_1^4h_0^2 = 0, \quad (19)$$

$$(G(\zeta))^3: 2anf_0f_1 + 2n^2c_5f_0^2f_1 + 4n^2c_6f_0^3f_1 + 4n^2c_2f_0^3f_1^3 + 10n^2c_3f_0^2f_1^3 + 20n^2c_4f_0f_1^3 + 2n^2c_5f_1^3h_0 + 20n^2c_6f_0f_1^3h_0 + n^2c_1f_1^3 = 0, \quad (20)$$

$$(G(\zeta))^5: 2n^2c_5f_1^3 + 16n^2c_6f_0f_1^3 + n^2c_3f_1^5 + 6n^2c_4f_0f_1^5 = 0 \quad (21)$$

$$(G(\zeta))^6: 6n^2c_6f_1^4 + n^2c_4f_1^6 = 0. \quad (22)$$

The following results are obtained *via* solving this system:

Result 1.

$$f_0 = -\frac{2a(2+n)h_0}{n^2(c_1 - 4c_5h_0)}, f_1 = -\frac{2a(2+n)\sqrt{-h_0}}{\sqrt{n^4(c_1 - 4c_5h_0)^2}},$$

$$c_4 = \frac{3n^4c_6(c_1 - 4c_5h_0)^2}{2a^2(2+n)^2h_0},$$

$$c_2 = \frac{\left(n^2(1+n)c_1^2 + 4n^2(4+n)c_1c_5h_0 + 16(-n^2(5+2n)c_5^2 + 4a(2+n)^2c_6)h_0^2 \right)}{4a(2+n)^2h_0},$$

$$c_3 = \frac{\left[n^2(c_1 - 4c_5h_0) \times (n^2c_1c_5 + 4(-n^2c_5^2 + 5a(2+n)c_6)h_0) \right]}{2a^2(2+n)^2h_0},$$

$$w = -ak^2 - \frac{4ah_0}{n^2}. \quad (23)$$

Result 2.

$$f_0 = \frac{\sqrt{2h_0(c_3c_5 + 50c_6^2h_0)} - 10c_6h_0}{c_3}, f_1 = \frac{\pm 1}{\sqrt{c_3}} \times$$

$$\times \sqrt{\frac{20c_6(\sqrt{2h_0(c_3c_5 + 50c_6^2h_0)} - 10c_6h_0)}{c_3} - 2c_5},$$

$$k = \frac{\sqrt{-n^2w - 4ah_0}}{\sqrt{an}},$$

$$c_1 = 4c_5h_0 - \frac{a(2+n) \left(\frac{10c_6h_0 + \sqrt{2h_0(c_3c_5 + 50c_6^2h_0)}}{\sqrt{2h_0(c_3c_5 + 50c_6^2h_0)}} \right)}{n^2c_5},$$

$$c_2 = \frac{a(1+n) \left(c_3c_5 + 10c_6 \left(\frac{10c_6h_0 + \sqrt{2h_0(c_3c_5 + 50c_6^2h_0)}}{\sqrt{2h_0(c_3c_5 + 50c_6^2h_0)}} \right) \right)}{2n^2c_5^2} - 14c_6h_0 - \sqrt{18h_0(c_3c_5 + 50c_6^2h_0)},$$

$$c_4 = \frac{3c_6 \left(c_3c_5 + 10c_6 \left(\frac{10c_6h_0 + \sqrt{2h_0(c_3c_5 + 50c_6^2h_0)}}{\sqrt{2h_0(c_3c_5 + 50c_6^2h_0)}} \right) \right)}{c_5^2}. \quad (24)$$

Result 3.

$$f_1 = \pm \frac{\sqrt{2a(5+2n)f_0 - 32n^2c_6f_0^3}}{\sqrt{n^2(3c_1 + 2c_2f_0)}},$$

$$w = \frac{a(-ak^2(5+2n) + 6c_1f_0 + 4(c_2 + 4k^2n^2c_6)f_0^2)}{a(5+2n) - 16n^2c_6f_0^2},$$

$$c_3 = \frac{\left(c_1(a + an - 44n^2c_6f_0^2) + 2c_2f_0(a(2+n) - 20n^2c_6f_0^2) \right)}{-2a(5+2n)f_0^2 + 32n^2c_6f_0^4},$$

$$c_4 = \frac{3n^2c_6(3c_1 + 2c_2f_0)}{16n^2c_6f_0^3 - a(5+2n)f_0},$$

$$c_5 = \frac{2a(2+n)c_2f_0 + c_1(a + an + 16n^2c_6f_0^2)}{2n^2f_0(3c_1 + 2c_2f_0)},$$

$$h_0 = -\frac{n^2f_0(3c_1 + 2c_2f_0)}{2a(5+2n) - 32n^2c_6f_0^2}. \quad (25)$$

Substituting equations (10) and (11) into equation (8) yields a polynomial expressed in terms of powers of $G(\zeta)$. We then arrange terms with similar powers and set each corresponding coefficient to zero. This process generates the following system of algebraic equations:

$$q_1(x, t) = e^{i \left(u - kx + \frac{t^\mu(-ak^2 - \frac{4ah_0}{n^2})}{\mu} \right)} \times$$

$$\left(-\frac{2a(2+n)h_0}{n^2(c_1 - 4c_5h_0)} \pm \eta_1 \tanh \left[\left(x + \frac{2akt^\mu}{\mu} \right) \sqrt{-h_0} \right] \right)^{\frac{1}{n}}, \quad (26)$$

where $\eta_1 = \frac{2a(2+n)\sqrt{-h_0}\sqrt{-h_0}}{\sqrt{n^4(c_1 - 4c_5h_0)^2}}$

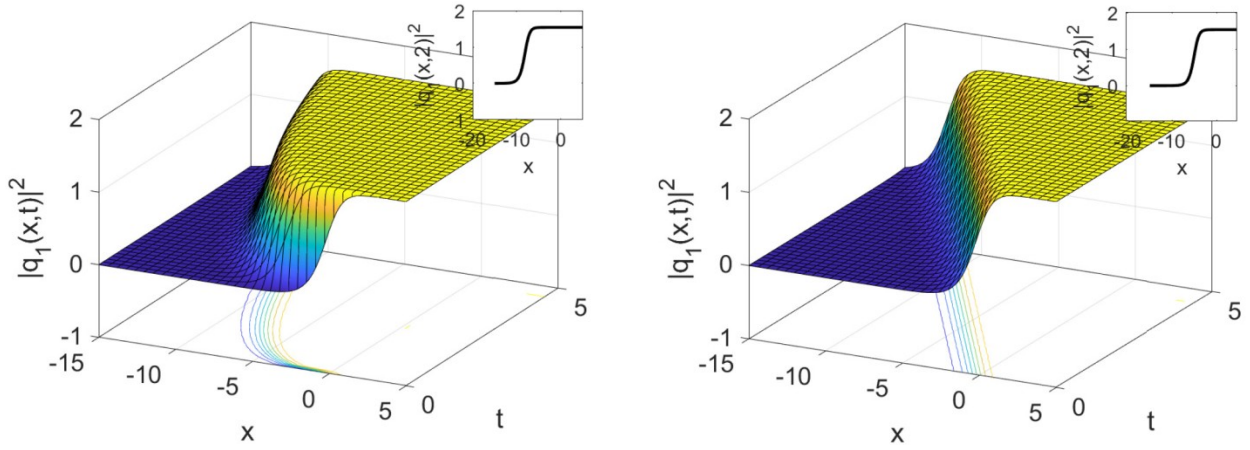


Fig. 1. The comparison of kink-type plots $|q_1(x,t)|^2$ for $\mu = 1$ and $\mu = 0.4$, where $n = 3$, $c_1 = 1$, $c_5 = 0.1$, $u = 0.1$, $a = 1.2$, $k = 1$, $h_0 = -1$, and $s = 1$.

$$q_2(x,t) = e^{i\left(u-kx + \frac{t^\mu(-ak^2 - \frac{4ah_0}{n^2})}{\mu}\right)} \times \left(\frac{-2a(2+n)h_0}{n^2(c_1-4c_5h_0)} \pm \eta_1 \coth\left[\left(s+x + \frac{2akt^\mu}{\mu}\right)\sqrt{-h_0}\right]\right)^{\frac{1}{n}}, \quad (27)$$

$$q_3(x,t) = e^{i\left(u-kx + \frac{t^\mu(-ak^2 - \frac{4ah_0}{n^2})}{\mu}\right)} \times \left(\eta_1 \left(\operatorname{isech}\left[2\sqrt{-h_0}\zeta\right] + \tanh\left[2\sqrt{-h_0}\zeta\right]\right)\right)^{\frac{1}{n}}, \quad (28)$$

where $\zeta = s + x + \frac{2akt^\mu}{\mu}$.

$$q_4(x,t) = e^{i\left(u-kx + \frac{t^\mu(-ak^2 - \frac{4ah_0}{n^2})}{\mu}\right)} \times \left(\eta_1 \left(\operatorname{csch}\left[2\sqrt{-h_0}\zeta\right] - \coth\left[2\sqrt{-h_0}\zeta\right]\right)\right)^{\frac{1}{n}}, \quad (29)$$

$$q_5(x,t) = e^{i\left(u-kx + \frac{t^\mu(-ak^2 - \frac{4ah_0}{n^2})}{\mu}\right)} \times \left(\eta_1 \left(\coth\left[\frac{\sqrt{-h_0}\zeta}{2}\right] + \tanh\left[\frac{\sqrt{-h_0}\zeta}{2}\right]\right)\right)^{\frac{1}{n}}, \quad (30)$$

Here, $\zeta = s + x + \frac{2\sqrt{at^\mu}\sqrt{-n^2w-4ah_0}}{n\mu}$

and

$$\eta_2 = \frac{\sqrt{-h_0}}{\sqrt{c_3}} \left(-2c_5 + \frac{20c_6(-10c_6h_0 + \sqrt{2}h_0(c_3c_5 + 50c_6^2h_0))}{c_3}\right)^{\frac{1}{2}}. \quad (31)$$

$$q_6(x,t) = e^{i\left(u + \frac{wt^\mu}{\mu} - x\sqrt{\frac{-n^2w-4ah_0}{\sqrt{an}}}\right)} \times \left(\frac{-10c_6h_0 + \sqrt{-2h_0(c_3c_5 + 50c_6^2h_0)}}{c_3} \pm \eta_2 \tanh\left[\sqrt{-h_0}\zeta\right]\right)^{\frac{1}{n}}, \quad (32)$$

$$q_7(x,t) = e^{i\left(u + \frac{wt^\mu}{\mu} - x\sqrt{\frac{-n^2w-4ah_0}{\sqrt{an}}}\right)} \times \left(\frac{-10c_6h_0 + \sqrt{-2h_0(c_3c_5 + 50c_6^2h_0)}}{c_3} \pm \eta_2 \coth\left[\sqrt{-h_0}\zeta\right]\right)^{\frac{1}{n}}, \quad (33)$$

$$q_8(x,t) = e^{i\left(u + \frac{wt^\mu}{\mu} - x\sqrt{\frac{-n^2w-4ah_0}{\sqrt{an}}}\right)} \times \left(f_0 \pm \eta_2 \left(\operatorname{isech}\left[2\sqrt{-h_0}\zeta\right] + \tanh\left[2\sqrt{-h_0}\zeta\right]\right)\right)^{\frac{1}{n}}, \quad (34)$$

$$q_9(x,t) = e^{i\left(u + \frac{wt^\mu}{\mu} - x\sqrt{\frac{-n^2w-4ah_0}{\sqrt{an}}}\right)} \times \left(f_0 \pm \eta_2 \left(-\coth\left[2\sqrt{-h_0}\zeta\right] + \operatorname{csch}\left[2\sqrt{-h_0}\zeta\right]\right)\right)^{\frac{1}{n}}, \quad (35)$$

$$q_{10}(x,t) = e^{i\left(u + \frac{wt^\mu}{\mu} - x\sqrt{\frac{-n^2w-4ah_0}{\sqrt{an}}}\right)} \times \left(f_0 \pm \frac{\eta_2}{2} \left(\coth\left[\frac{\sqrt{-h_0}}{2}\zeta\right] + \tanh\left[\frac{\sqrt{-h_0}}{2}\zeta\right]\right)\right)^{\frac{1}{n}}, \quad (36)$$

$$q_{11}(x,t) = e^{i\left(u-kx + \frac{at^\mu(-ak^2(5+2n) + 6c_1f_0 + 4(c_2+4k^2n^2c_6)f_0^2)}{\mu(a(5+2n)-16n^2c_6f_0^2)}\right)} \times \left(f_0 \pm \eta_3 \tanh[H_1]\right)^{\frac{1}{n}}, \quad (37)$$

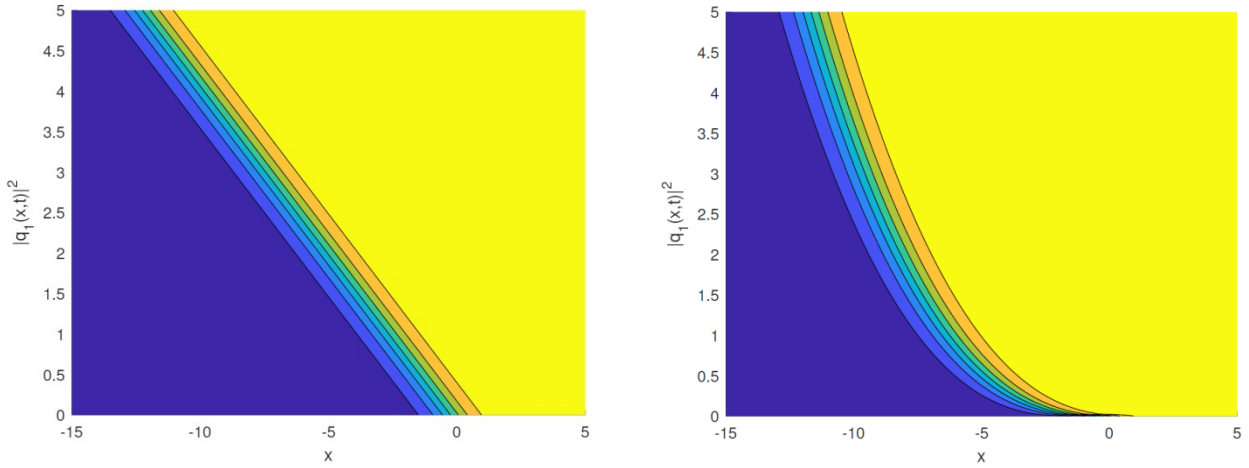


Fig. 2. The comparison of contour plots $Re(q_1(x,t))$ for $\mu = 1$ and $\mu = 0.4$, where $n = 3, c_1 = 1, c_5 = 0.1, u = 0.1, a = 1.2, k = 1, h_0 = -1$, and $s = 1$.

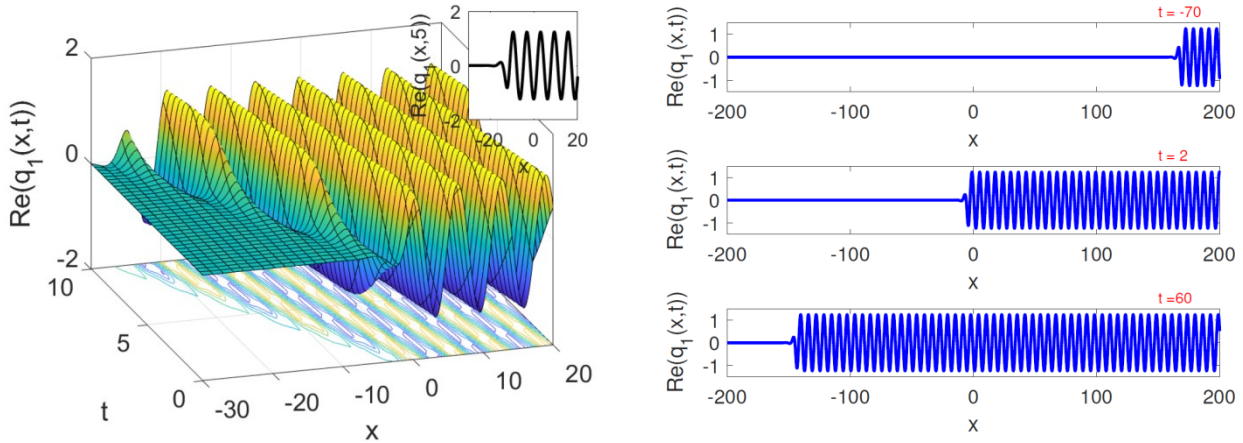


Fig. 3. The wave plots of $|q_1(x,t)|^2$ with the effect of temporal time parameter, where $n = 3, c_1 = 1, c_5 = 0.1, u = 0.1, a = 1.2, k = 1, h_0 = -1, \mu = 1$, and $s = 1$.

where $H_1 = \left(s + x + \frac{2akt^\mu}{\mu} \right) \sqrt{\frac{n^2 f_0 (3c_1 + 2c_2 f_0)}{2a(5+2n) - 32n^2 c_6 f_0^2}}$ and

$$\eta_3 = \sqrt{\frac{n^2 f_0 (3c_1 + 2c_2 f_0) (2a(5+2n) f_0 - 32n^2 c_6 f_0^3)}{(2a(5+2n) - 32n^2 c_6 f_0^2) (n^2 (3c_1 + 2c_2 f_0))}}$$

$$q_{12}(x,t) = e^{i\left(u - kx + \frac{at^\mu(-ak^2(5+2n) + 6c_1 f_0 + 4(c_2 + 4k^2 n^2 c_6) f_0^2)}{\mu(a(5+2n) - 16n^2 c_6 f_0^2)}\right)} (f_0 \pm \eta_3 \coth[H_1])^{\frac{1}{n}}, \quad (38)$$

$$q_{13}(x,t) = e^{i\left(u - kx + \frac{at^\mu(-ak^2(5+2n) + 6c_1 f_0 + 4(c_2 + 4k^2 n^2 c_6) f_0^2)}{\mu(a(5+2n) - 16n^2 c_6 f_0^2)}\right)} (f_0 \pm \eta_3 (\operatorname{isech}[2H_1] + \tanh[2H_1]))^{\frac{1}{n}}, \quad (39)$$

$$q_{14}(x,t) = e^{i\left(u - kx + \frac{at^\mu(-ak^2(5+2n) + 6c_1 f_0 + 4(c_2 + 4k^2 n^2 c_6) f_0^2)}{\mu(a(5+2n) - 16n^2 c_6 f_0^2)}\right)} (f_0 \pm \eta_3 (-\coth[2H_1] + \operatorname{csch}[2H_1]))^{\frac{1}{n}}, \quad (40)$$

$$q_{15}(x,t) = e^{i\left(u - kx + \frac{at^\mu(-ak^2(5+2n) + 6c_1 f_0 + 4(c_2 + 4k^2 n^2 c_6) f_0^2)}{\mu(a(5+2n) - 16n^2 c_6 f_0^2)}\right)} \left(f_0 \pm \eta_3 \left(\coth\left[\frac{H_1}{2}\right] + \tanh\left[\frac{H_1}{2}\right] \right) \right)^{\frac{1}{n}}. \quad (41)$$

For completeness, we provide graphical simulations of the obtained solutions. The solution (26) is illustrated in Figs. 1–3 and 8 (kink-type, contour, wave and bright-profile comparisons). The solution (32) is displayed in Figs. 4 and 5 (wave and contour plots).

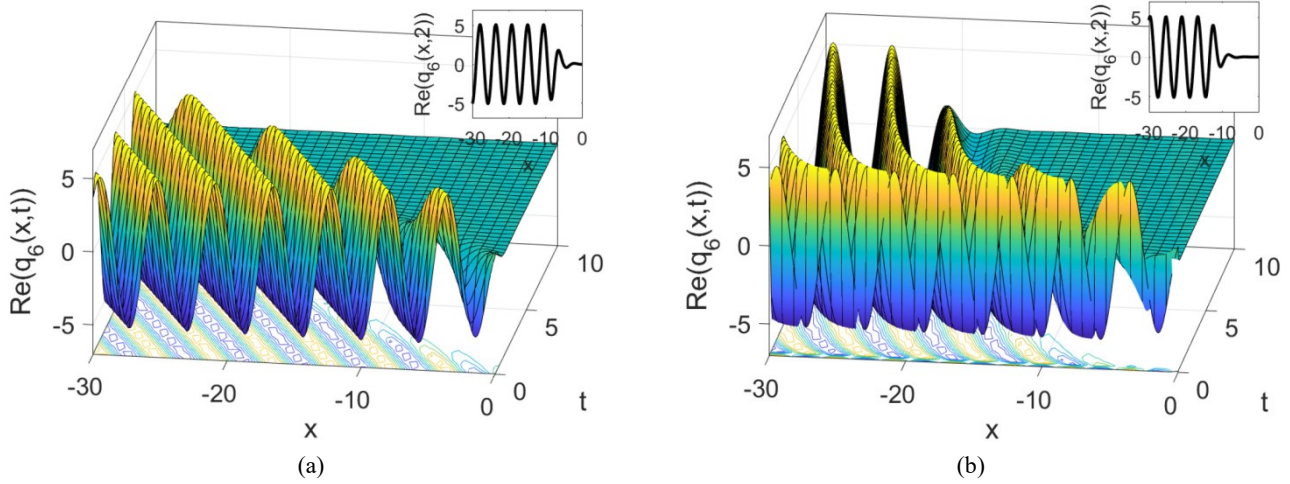


Fig. 4. The comparison of wave plots of $Re(q_6(x,t))$ for $\mu = 1$ and $\mu = 0.4$, where $n = 3, c_3 = -0.3, c_5 = 0.1, c_5 = 1, u = 0.1, a = 1.2, w = -2, h_0 = -1, \mu = 1$, and $s = 1$.

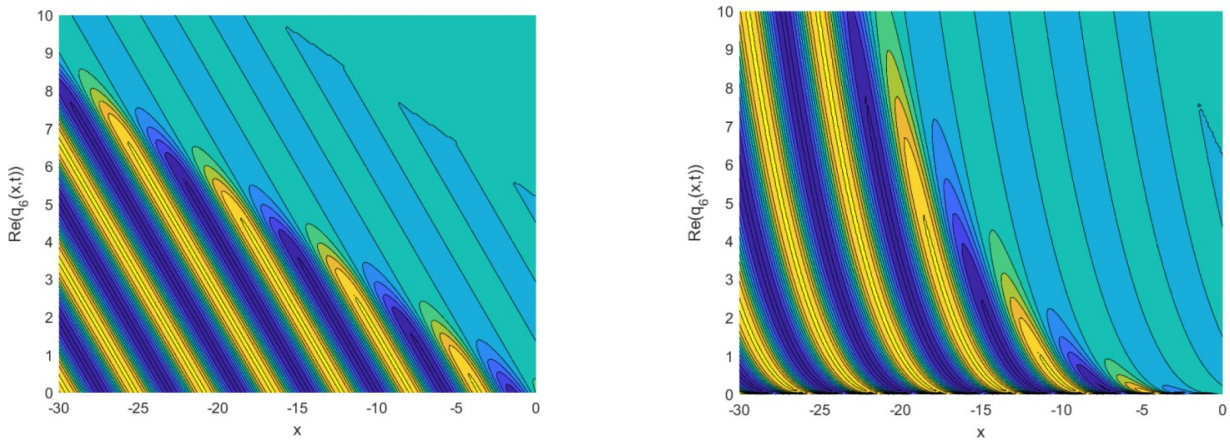


Fig. 5. The comparison of contour plots of $Re(q_6(x,t))$ for $\mu = 1$ and $\mu = 0.4$, where $n = 3, c_3 = -0.3, c_5 = 0.1, c_5 = 1, u = 0.1, a = 1.2, w = -2, h_0 = -1, \mu = 1$, and $s = 1$.

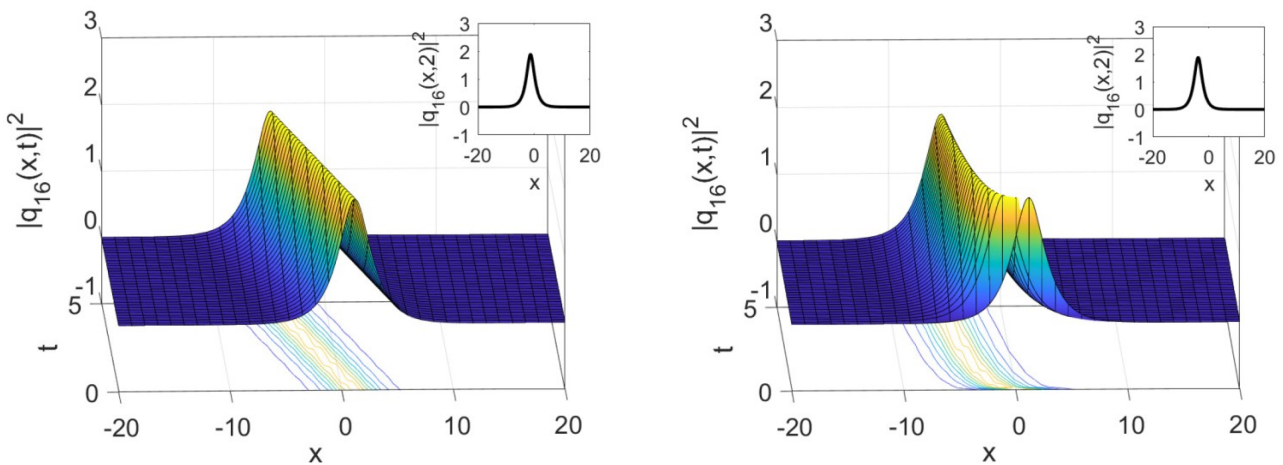


Fig. 6. The comparison of bright plots of $|q_{16}(x,t)|^2$ with the effect of different temporal parameter, where $c_5 = 1, n = 3, u = 0.1, a = 2, k = 0.3, c_3 = 0.3, s = 0.2, A = 0.1, B = 1$.

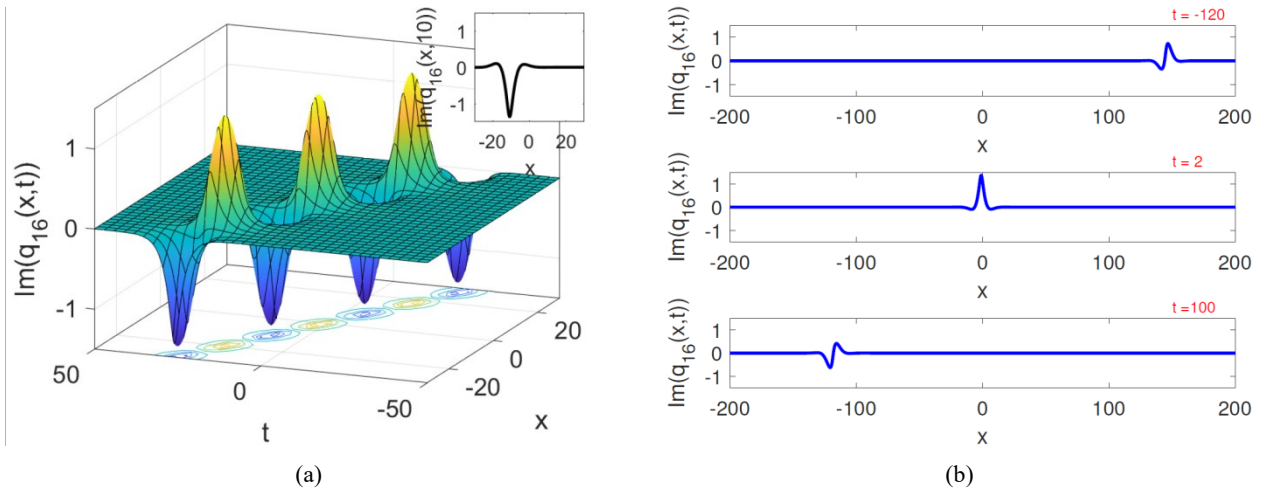


Fig. 7. The dark-bright plots of $Im(q_{16}(x,t))$ with the effect of the temporal parameter, where $c_5 = 1, n = 3, u = 0.1, a = 2, k = 0.3, c_3 = 0.3, s = 0.2, \mu = 1, A = 0.1, B = 1$.

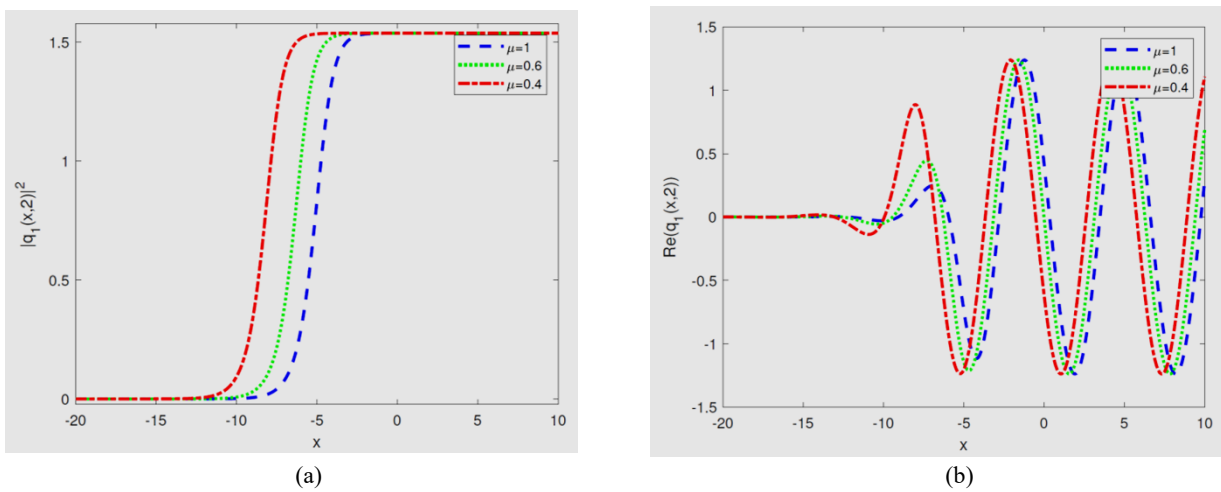


Fig. 8. The comparison of bright plots $|q_1(x,t)|^2$ and $Re(q_1(x,t))$ for $\alpha = 1$ and $\alpha = 0.4$, where $n = 3, c_1 = 1, c_5 = 0.1, u = 0.1, a = 1.2, k = 1, h_0 = -1$, and $s = 1$.

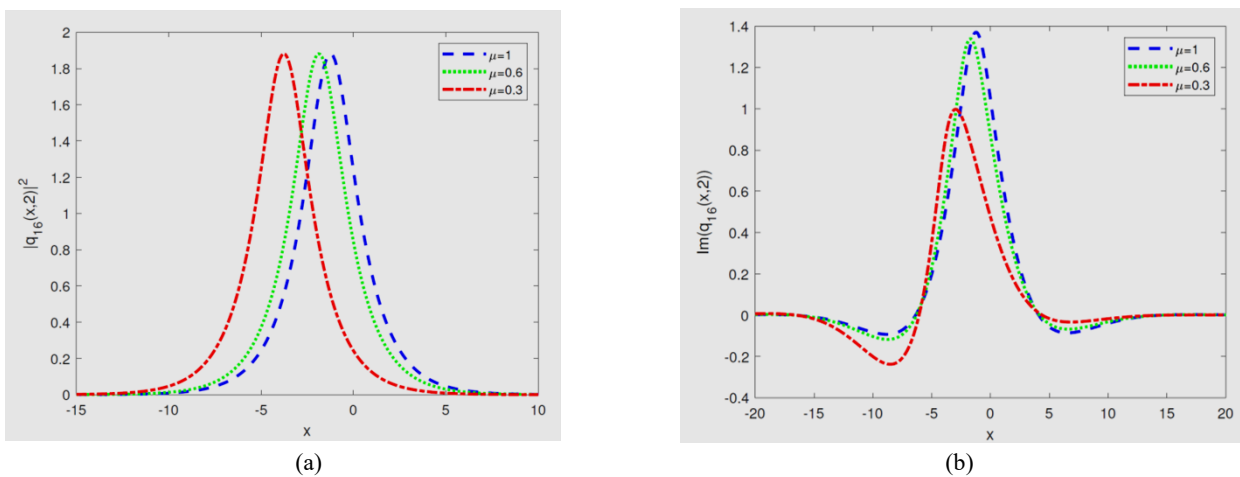


Fig. 9. The effect of the conformable parameter on $|u_{17}(x,t)|^2$ and $Re(u_{17}(x,t))$, where $c_5 = 1, n = 3, u = 0.1, a = 2, k = 0.3, c_3 = 0.3, s = 0.2, A = 0.1, B = 1$.

4. Kudryashov’s approach

In this section, several new conformable optical soliton solutions to the current model are constructed using the Kudryashov method. The method relies on the following function [28]:

$$K(\zeta) = \frac{1}{(A-B)\sinh(\zeta)+(A+B)\cosh(\zeta)}, \tag{42}$$

where A and B are non-zero constants, and the functions $K(\zeta)$ satisfies the following relation:

$$(K'(\zeta))^2 - K^2(\zeta)(1 - 4ABK(\zeta)^2) = 0. \tag{43}$$

According to the present method, the following series, based on the base function (42), is the solution of the proposed model:

$$V(\zeta) = \sum_{i=0}^N f_i(K(\zeta))^i, f_N \neq 0. \tag{44}$$

where the values f_0, f_1, f_2, \dots are the constants that need to be found, and N is the balancing value. Hence, the equation (44) has the following form:

$$V(\zeta) = f_0 + f_1 G(\zeta). \tag{45}$$

By substituting equation (45) and its derivatives, together with equation (3), into equation (8), we obtain a polynomial expression for $K(\zeta)$. After gathering all the coefficients of $K(\zeta)$ and equating them to zero, we form a system of algebraic equations. Solving this system yields the following results:

Result 1.

$$f_0 = 0, f_1 = \pm \frac{\sqrt{8ABC_5}}{\sqrt{c_3}}, w = a \left(\frac{1}{n^2} - k^2 \right), \tag{46}$$

$$c_1 = -c_5, c_2 = \frac{a(1+n)c_3}{2n^2c_5} - \frac{4c_4c_5}{3c_3}, c_6 = \frac{c_4c_5}{3c_3}.$$

Using equations (3), (6), (7), (42), (45), and (46), we can have the following optical solutions:

$$q_{16}(x, t) = e^{i \left(u - kx + \frac{a \left(\frac{1}{n^2} - k^2 \right) t^\mu}{\mu} \right)} \times \left(- \frac{\sqrt[3]{2^2 \sqrt{ABC_5}}}{((A+B)\cosh[\zeta] + (A-B)\sinh[\zeta])\sqrt{c_3}} \right)^{\frac{1}{n}}, \tag{47}$$

where $\zeta = s + x + \frac{2akt^\mu}{\mu}$.

Plugging $A = B$ into the above solution, we obtain the following bright soliton solution

$$q_{17}(x, t) = e^{i \left(u - kx + \frac{a \left(-k^2 + \frac{1}{n^2} \right) t^\mu}{\mu} \right)} \times \left(\pm \frac{\operatorname{sech} \left[s + x + \frac{2akt^\mu}{\mu} \right] \sqrt{2c_5}}{\sqrt{c_3}} \right)^{\frac{1}{n}}. \tag{48}$$

Plugging $A = -B$ into the above solution, we obtain the following soliton solution

$$q_{18}(x, t) = e^{i \left(u - kx + \frac{a \left(-k^2 + \frac{1}{n^2} \right) t^\mu}{\mu} \right)} \times \left(\pm \frac{\sqrt{-2B} \operatorname{csch} \left[s + x + \frac{2akt^\mu}{\mu} \right] \sqrt{c_5}}{\sqrt{c_3 B}} \right)^{\frac{1}{n}}. \tag{49}$$

Result 2.

$$f_0 = 0, f_1 = \pm \frac{\sqrt{8ABC_5}}{\sqrt{c_3}}, k = -\frac{\sqrt{a-n^2w}}{\sqrt{an}}, \tag{50}$$

$$c_1 = -c_5, c_4 = \frac{3c_3(a(1+n)c_3 - 2n^2c_2c_5)}{8n^2c_5^2},$$

$$c_6 = \frac{1}{8} \left(-2c_2 + \frac{a(1+n)c_3}{n^2c_5} \right).$$

Using equations (3), (6), (7), (42), (45), and (50), we can have the following optical solutions:

$$q_{19}(x, t) = e^{i \left(u + \frac{\sqrt{a-n^2wx} + t^\mu w}{\sqrt{an}} \right)} \times \left(\pm \frac{\sqrt[3]{2^2 \sqrt{ABC_5}}}{((A+B)\cosh[\zeta] + (A-B)\sinh[\zeta])\sqrt{c_3}} \right)^{\frac{1}{n}}, \tag{51}$$

where $\zeta = s + x - \frac{2\sqrt{at^\mu}\sqrt{a-n^2w}}{n\mu}$.

Plugging $A = B$ into the above solution, we obtain the following bright soliton solution

$$q_{20}(x, t) = e^{i \left(u + \frac{\sqrt{a-n^2wx} + t^\mu w}{\sqrt{an}} \right)} \times \left(\pm \frac{\sqrt{2c_5}}{\sqrt{c_3}} \operatorname{sech} \left[s + x - \frac{2\sqrt{at^\mu}\sqrt{a-n^2w}}{n\mu} \right] \right)^{\frac{1}{n}}. \tag{52}$$

Plugging $A = -B$ into the above solution, we obtain the following soliton solution

$$q_{21}(x, t) = e^{i \left(u + \frac{\sqrt{a-n^2wx} + t^\mu w}{\sqrt{an}} \right)} \times \left(\pm \frac{\sqrt{-2Bc_5}}{\sqrt{Bc_3}} \operatorname{csch} \left[s + x - \frac{2\sqrt{at^\mu}\sqrt{a-n^2w}}{n\mu} \right] \right)^{\frac{1}{n}}. \tag{53}$$

Result 3.

$$f_0 = 0, f_1 = \pm \frac{\sqrt{-8ABc_1}}{\sqrt{c_3}}, k = -\frac{\sqrt{a-n^2w}}{\sqrt{an}}, \tag{54}$$

$$c_5 = -c_1, c_4 = \frac{3c_3(2n^2c_1c_2 + a(1+n)c_3)}{8n^2c_1^2},$$

$$c_6 = -\frac{c_2}{4} - \frac{a(1+n)c_3}{8n^2c_1}.$$

Using equations (3), (6), (7), (42), (45), and (54), we can have the following optical solutions:

$$q_{22}(x, t) = e^{i \left(u + \frac{\sqrt{a-n^2wx} + t^\mu w}{\sqrt{an}} \right)} \times \left(\pm \frac{\sqrt[3]{-2^2 \sqrt{ABc_1}}}{((A+B)\cosh[\zeta] + (A-B)\sinh[\zeta])\sqrt{c_3}} \right)^{\frac{1}{n}}, \tag{55}$$

where $\zeta = s + x - \frac{2\sqrt{at^\mu}\sqrt{a-n^2w}}{n\mu}$.

Plugging $A = B$ into the above solution, we obtain the following bright soliton solution

$$q_{23}(x, t) = e^{i\left(u + \frac{\sqrt{a-n^2}wx + t^\mu w}{\sqrt{an}}\right)} \times \left(\pm \frac{\sqrt{-2c_1}}{\sqrt{c_3}} \operatorname{sech}\left[x - \frac{2\sqrt{a}t^\mu\sqrt{a-n^2}w}{n\mu}\right]\right)^{\frac{1}{n}}. \quad (56)$$

Plugging $A = -B$ into the above solution, we obtain the following soliton solution

$$q_{24}(x, t) = e^{i\left(u + \frac{\sqrt{a-n^2}wx + t^\mu w}{\sqrt{an}}\right)} \times \left(\frac{\sqrt{2Bc_1}}{\sqrt{Bc_3}} \operatorname{csch}\left[x - \frac{2\sqrt{a}t^\mu\sqrt{a-n^2}w}{n\mu}\right]\right)^{\frac{1}{n}}. \quad (57)$$

The solution (47) is presented in Figs. 6 and 7 (bright and dark-bright profiles). Finally, Fig. 9 demonstrates the influence of the conformable parameter on solution (48).

6. Conclusions

This paper derived the optical soliton solutions to the governing NLSE with fractional temporal evolution, linear CD, and having Kudryashov's form of SPM structure. Two approaches have made this retrieval possible. They are the Kudryashov integration scheme and the modified simplest equation method. These two approaches collectively yielded the soliton solutions, which are enumerated and listed along with their respective parameter constraints that guarantee the existence of the solitons. The results are thus applicable to the telecommunications industry, which is the lifeline of modern-day internet communication, keeping the entire planet connected.

The future of this study stands on a very strong footing. The results will be later studied for additional form of fibers, namely polarization-preserving fibers, dispersion-flattened fibers, dispersion-managed fibers, etc. The model can also be applied to several forms of optoelectronic devices such as optical metamaterials and metasurfaces, optical couplers, magneto-optic waveguides, Bragg gratings, and others. The results of these research activities will be disseminated once they start pouring in sequentially.

Disclosures

Competing interests

The authors declare that they have no competing interests.

Acknowledgement

The work for the fourth author (AB) is supported by Grambling State University for the Endowed Chair of Mathematics and this support is gratefully acknowledged.

References

1. Biswas A., Pati G. S. Mathematical theory of slow light optical solitons. *Waves in Random and*

- Complex Media*. 2011. **21**, No 3. P. 456–468. <https://doi.org/10.1080/17455030.2011.582892>.
2. González-Gaxiola O., Biswas A. Belic M.R. Optical soliton perturbation of Fokas–Lenells equation by the Laplace–Adomian decomposition algorithm. *J. Eur. Opt. Soc. – Rapid Publ.* 2019. **15**. P. 13. <https://doi.org/10.1186/s41476-019-0111-6>.
3. Murad M.A.S., Arnous A.H., Biswas A. *et al.* Suppressing internet bottleneck with Kudryashov's extended version of self-phase modulation and fractional temporal evolution. *J. Opt.* 2024. <https://doi.org/10.1007/s12596-024-01937-4>.
4. Biswas A., Jawad A.J.M., Manrakhani W.N. *et al.* Optical solitons and complexitons of the Schrödinger–Hirota equation. *Opt. Laser Technol.* 2012. **44**, No 7. P. 2265–2269. <https://doi.org/10.1016/j.optlastec.2012.02.028>.
5. Kudryashov N.A. Method for finding optical solitons of generalized nonlinear Schrödinger equations. *Optik*. 2022. **261**. P. 169163. <https://doi.org/10.1016/j.ijleo.2022.169163>.
6. Kudryashov N.A. Method for finding highly dispersive optical solitons of nonlinear differential equations. *Optik*. 2020. **206**. P. 163550. <https://doi.org/10.1016/j.ijleo.2019.163550>.
7. Murad M.A.S. Optical solutions for perturbed conformable Fokas–Lenells equation via Kudryashov auxiliary equation method. *Mod. Phys. Lett. B*. 2025. **39**, No 07. P. 2450418. <https://doi.org/10.1142/S0217984924504189>.
8. Ali K., Yusuf A., Alquran M., Tarla S. New physical structures and patterns to the optical solutions of the nonlinear Schrödinger equation with a higher dimension. *Commun. Theor. Phys.* 2023. **75**. P. 085003. <https://doi.org/10.1088/1572-9494/acde69>.
9. Murad M.A.S., Faridi W.A., Iqbal M. *et al.* Analysis of Kudryashov's equation with conformable derivative via the modified Sardar sub-equation algorithm. *Results Phys.* 2024. **60**. P. 107678. <https://doi.org/10.1016/j.rinp.2024.107678>.
10. Murad M.A.S. Formation of optical soliton wave profiles of nonlinear conformable Schrödinger equation in weakly non-local media: Kudryashov auxiliary equation method. *J. Opt.* 2025. **54**. P. 3177–3190. <https://doi.org/10.1007/s12596-024-02110-7>.
11. Mahmud F., Samsuzzoha M., Akbar M.A. The generalized Kudryashov method to obtain exact traveling wave solutions of the PHI-four equation and the Fisher equation. *Results Phys.* 2017. **7**. P. 4296–4302. <https://doi.org/10.1016/j.rinp.2017.10.049>.
12. Murad M.A.S. Perturbation of optical solutions and conservation laws in the presence of a dual form of generalized nonlocal nonlinearity and Kudryashov's refractive index having quadrupled power-law. *Opt. Quantum Electron.* 2024. **56**, No 5. P. 864. <https://doi.org/10.1007/s11082-024-06676-1>.
13. Wang K.-J. Multi-wave complexiton, multi-wave, interaction-wave and the travelling wave solutions to the (2+1)-dimensional Boiti–Leon–Manna–Pempinelli equation for the incompressible fluid. *Pramana*. 2024. **98**. P. 47. <https://doi.org/10.1007/s12043-024-02725-2>.

14. Kudryashov N.A. A generalized model for description of propagation pulses in optical fiber. *Optik*. 2019. **189**. P. 42–52. <https://doi.org/10.1016/j.ijleo.2019.05.069>.
15. Mathanaranjan T., Kumar D., Rezazadeh H., Akinyemi L. Optical solitons in metamaterials with third and fourth order dispersions. *Opt. Quant. Electron*. 2022. **54**, No 5. P. 271. <https://doi.org/10.1007/s11082-022-03656-1>.
16. Iqbal M., Seadawy A.R., Lu D., Zhang Z. Physical structure and multiple solitary wave solutions for the nonlinear Jaulent–Miodek hierarchy equation. *Mod. Phys. Lett. B*. 2024. **38**, No. 16. P. 2341016. <https://doi.org/10.1142/S0217984923410166>.
17. Faridi W.A., Bakar M.A., Akgül A. *et al.* Exact fractional soliton solutions of thin-film ferroelectric material equation by analytical approaches. *Alexandria Eng. J.* 2023. **78**. P. 483–497. <https://doi.org/10.1016/j.aej.2023.07.049>.
18. Mathanaranjan T. Optical soliton, linear stability analysis and conservation laws via multipliers to the integrable Kuralay equation. *Optik*. 2023. **290**. P. 171266. <https://doi.org/10.1016/j.ijleo.2023.171266>.
19. Elsherbeny A.M., El-Barkouky R., Ahmed H.M. *et al.* Optical soliton perturbation with Kudryashov’s generalized nonlinear refractive index. *Optik*. 2021. **240**. P. 166620. <https://doi.org/10.1016/j.ijleo.2021.166620>.
20. Yildirim Y., Biswas A., Ekici M. *et al.* Optical soliton perturbation with Kudryashov’s law of arbitrary refractive index. *J. Opt.* 2021. **50**. P. 245–252. <https://doi.org/10.1007/s12596-021-00693-z>.
21. Zayed E.M.E., Shohib R.M.A., Alngar M.E.M. *et al.* Optical solitons and conservation laws associated with Kudryashov’s sextic power-law nonlinearity of refractive index. *Ukr. J. Phys. Opt.* 2021. **22**, No 1. P. 38–49. <https://doi.org/10.3116/16091833/22/1/38/2021>.
22. Xu X.-Z. Exact chirped solutions for the NLSE having Kudryashov’s law with dual form of generalized non-local nonlinearity. *Optik*. 2023. **287**. P. 171101. <https://doi.org/10.1016/j.ijleo.2023.171101>.
23. Yildirim Y., Biswas A., Triki H. *et al.* Cubic–quartic optical soliton perturbation with Kudryashov’s law of refractive index having quadrupled–power law and dual form of generalized nonlocal nonlinearity by sine-Gordon equation approach. *J. Opt.* 2021. **50**. P. 593–599. <https://doi.org/10.1007/s12596-021-00686-y>.
24. Zayed E.M.E., Shohib R.M.A., Alngar M.E.M. Cubic–quartic optical solitons with Kudryashov’s arbitrary form of nonlinear refractive index. *Optik*. 2021. **238**. P. 166747. <https://doi.org/10.1016/j.ijleo.2021.166747>.
25. Gepreel K.A., Zayed E.M.E., Alngar M.E.M. *et al.* Optical solitons with Kudryashov’s arbitrary form of refractive index and generalized non-local nonlinearity. *Optik*. 2021. **243**. P. 166723. <https://doi.org/10.1016/j.ijleo.2021.166723>.
26. Khalil R., Al Horani M., Yousef A., Sababheh M. A new definition of fractional derivative. *J. Comput. Appl. Math.* 2014. **264**. P. 65–70. <https://doi.org/10.1016/j.cam.2014.01.002>.
27. Awadalla M., Zafar A., Taishiyeva A. *et al.* The analytical solutions to the M-fractional Kairat-II and Kairat-X equations. *Front. Phys.* 2023. **11**. P. 1335423. <https://doi.org/10.13140/RG.2.2.22148.30088>.
28. Ozisik M., Secer A., Bayram M., Aydin H. An encyclopedia of Kudryashov’s integrability approaches applicable to optoelectronic devices. *Optik*. 2022. **265**. P. 169499. <https://doi.org/10.1016/j.ijleo.2022.169499>.

Authors and CV



Muhammad Amin S. Murad is an Assistant Professor in the Department of Mathematics, College of Science, Duhok University, Iraq. He earned a master’s degree in Mathematics and Computing for Finance from Swansea University, Swansea, UK, in 2014, and PhD in Applied Mathematics from the University of Duhok, Duhok, Iraq, in 2023. His research interests include applied mathematics, fluid mechanics, numerical analysis and its applications in ordinary and partial differential equations. E-mail: muhammad.murad@uod.ac, <https://orcid.org/0000-0002-3402-1796>



Faraidun Kadir Hamasalh is a distinguished Professor of Mathematics specializing in Numerical Analysis. He is affiliated with the Sulaimani Polytechnic University and Department of Mathematics at the College of Education, University of Sulaimani, analysis and analytic methods. Iraq. Professor Hamasalh earned his PhD in Mathematics (Numerical Analysis) from the Department of Mathematics, College of Science, University of Sulaimani. His academic journey has been marked by a commitment to advancing research and teaching in mathematics, particularly in numerical analysis and analytic methods. <https://doi.org/0000-0001-8437-9043>



Ahmed H. Arnous is an Associate Professor in the Department of Engineering Mathematics and Physics at El-Shorouk Academy, Cairo, Egypt. With a strong academic foundation, he obtained his Master’s degree from Zagazig University, followed by PhD from Al-Azhar University. His research is deeply rooted in soliton theory and mathematical physics, focusing on their diverse applications. His contributions to the field are well recognized, making significant strides in advancing the understanding and application of mathematical concepts in physics. E-mail: ahmed.h.arnous@gmail.com, <https://orcid.org/0000-0002-7699-7068>



Anjan Biswas, PhD in Applied Mathematics (University of New Mexico in Albuquerque, NM, USA), post-doctoral researcher (University of Colorado, Boulder, CO, USA). Currently, he is the Endowed Chair of Mathematics at the Grambling

State University in Grambling, LA, USA. His current research interest is in quantum optoelectronics. His Erdos number is 4, while his H-index stands at 114. <http://orcid.org/0000-0002-8131-6044>



Anwar Jawad earned his PhD degree in Applied Mathematics from the University of Technology, Iraq in 2000. Currently, he is working as faculty staff at the Al Rafidain University College, Baghdad, Iraq. He is Professor of Applied Mathematics since 2014.

He is the author of more than 135 papers his h-index is 25, according to Google Scholar, one of his papers attracted more than 580 citations. His research areas include applied mathematics, statistics, and operation research. E-mail: anwar_jawad2001@yahoo.com, <https://orcid.org/0000-0001-8303-3235>



Yakup Yildirim, PhD in Mathematics (Uludag University, Turkey), Assistant Professor at the Biruni University, Turkey. The area of his scientific interests includes optical soliton solutions, conservation laws and Lie symmetry analysis.

He is the author of more than 300 publications. E-mail: yakupyildirim110@gmail.com, <http://orcid.org/0000-0003-4443-3337>



Luminita Moraru is currently Full Professor in the Department of Physics, Chemistry and Environment at Dunarea de Jos University of Galati, Romania. Her main areas of interest are image processing, pattern recognition, AI&ML&DL, modeling

and simulation. Also, her passion for nonlinear fiber optics, solitons is the secondary area of interest. She is author and co-author of more than 250 publications in refereed journals, national and international conferences, and book chapters. E-mail: lmoraru03@gmail.com, <https://orcid.org/0000-0002-9121-5714>



Carmelia Mariana Dragomir Balanica completed her PhD in Industrial Engineering in 2011 at the University, "Dunărea de Jos" in Galati, Romania, and her postdoctoral studies in Engineering Sciences were completed in 2015. She is an associate

professor at "Dunărea de Jos" in Galati, Romania. Her scientific interests include statistics and engineering sciences. <https://doi.org/0000-0001-7743-928X>

Authors' contributions

Murad M.A.S.: formal analysis, resources, data curation.

Hamasalh F.K.: formal analysis and resources.

Arnous A.H.: analysis, project administration, writing and review.

Biswas A.: project administration.

Jawad A.J.M.: writing, review and editing.

Yildirim Y.: writing, review and editing.

Moraru L.: formal analysis and project administration.

Dragomir C.: project administration and formal analysis.

Оптичні солітони форми Кудряшова із структурою локальної та нелокальної фазової само модуляції з дробовою часовою еволюцією

M.A.S. Murad, F.K. Hamasalh, A.H. Arnous, A. Biswas, A.J.M. Jawad, Y. Yildirim, L. Moraru, C. Dragomir

Анотація. У цій статті відновлено оптичні солітонні розв'язки за допомогою запропонованої Кудряшовим структури фазової само модуляції з дробовою часовою еволюцією. Ця модель може контролювати та пом'якшувати ефект вузького місця в мережі Інтернет, який є серйозною перешкодою для потоку інтернет-трафіку на міжконтинентальних відстанях. Два незалежні підходи до інтеграції зробили це відновлення можливим. Широкий спектр солітонних розв'язків з'явився завдяки колективному застосуванню схем інтеграції. Також представлено параметричні обмеження для існування таких солітонів.

Ключові слова: солітони, вузьке місце в мережі Інтернет, інтегрованість.