Plasmon-enhanced fluorometry based on gold nanostructure arrays. Method and device

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Abstract. In this work, we describe a method of surface-enhanced fluorometry, based on the phenomenon of localized surface plasmon resonance in unordered gold nanostructure arrays. The theoretical approach for the model system “gold nanoparticle-dielectric spacer” in the electrostatic approximation by solution of Laplace’s equation is considered. The developed technology for manufacturing the plasmonic substrates as well as design of the novel laser-based compact fluorometer are presented. The arrays of gold nanostructures on solid substrates (nanochips) coated with different thicknesses of SiO2 were developed and fabricated by thermal annealing of gold island films with subsequent dielectric spacer deposition. As an example for verification of the proposed method, the fluorescence properties of the system “gold nanostructures array – SiO2 dielectric coating – Rhodamine 6G” were studied. It has been shown that enhancement of dye emission up to 22 times for dielectric coating with the thickness of about 20 nm is possible. Presented method is of importance for the development of the novel nanoscale sensors, biomolecular assays and nanoplasmonic devices.

Keywords: localized surface plasmon resonance, plasmonic nanostructures, plasmon-enhanced fluorescence, rhodamine 6G.

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1. Introduction

Today, the number of works related to the investigation of fluorescence, one of the most subtle and complicated optical phenomena, is increasing rapidly. This is caused by the prospects for using fluorescence as a basis of rapid methods for detection and identification of chemical and biological substances. Fluorescence is one of the most widely used methods in biochemistry [1], clinical testing [2] and cellular imaging [3]. The capabilities of fluorescence technology are now exploited for quantitative and qualitative analysis and for biological research in general [4]. For example, structure of cells and their organelles, ion fluxes in cells and membrane processes have been investigated with application of fluorescent probes [5]. Additionally, fluorescence methods are being increasingly used for studies of conformational changes of biomolecules and specific interactions between them [6]. Fluorescence technique can also be an effective tool in development of new medical treatments [7] and drug delivery methods [8]. A large number of studies using fluorescence methods in biochemistry and molecular biology have been focused on using dyes to visualize individual components of biological systems [9, 10].

Many researches reported on the phenomenon of plasmon-enhanced fluorescence (PEF), which is induced by localized surface plasmon resonance (LSPR) that appears in high-conductive metal (plasmonic) nanostructures (PN) [11] and can significantly increase the emission of fluorophore [12, 13]. In particular, a large change in light emission of fluorescent molecule

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located near PN is caused by high electromagnetic field around gold and silver PN under LSPR conditions [14]. A number of active structures have been developed for PEF, primarily based on silver colloids [15, 16], silver and gold island films [17, 18], and ordered noble metal nanoparticle arrays [19].

In order to observe PEF, fluorescent molecule should be located near to the LSPR-generating surface, i.e. in the PN near-field zone [20]. However, very close disposition of fluorophore and PN leads to fluorescence quenching caused by the energy transfer from the excited molecule to the metallic surface [21]. Common approach reported in literature to overcome this issue is to hold the fluorophore away from a nanostructured surface at an optimal distance by a dielectric coating, which provides stable PEF conditions [22]. Another approach depends on preparation of the complex: PN – chemical linker – biomolecule – dye, which provides the required distance from dye to nanoparticle [23]. These techniques make it possible to obtain up to 50-fold fluorescence enhancement [24].

In this paper, we describe the method for plasmon-enhanced fluorometry and technology for producing gold nanochips aimed at in-practice realization of the enhancement mechanism. Presented here are the experimental results of studying the enhanced fluorescence signal of organic dye Rhodamine 6G (R6G) located at different distances from gold nanostructures, which was realized by using SiO₂ dielectric coatings. We also report on the application of developed and patented by authors [25] compact laser-based fluorometer “FluorotestNano” to detect the enhanced fluorescence emission of the dye using the above-mentioned method. We believe that presented methodology will contribute to the ongoing search for the combination of fluorescence spectroscopy methods, nanoplasmonics and nanotechnology in order to establish the possibilities of plasmon-enhanced fluorescence towards building highly efficient PEF-sensors.

2. Theory and modeling

The relationship between the plasmon-induced electromagnetic near-field of a metal nanostructure and properties of the adjacent dye has been previously studied in [26]. The fluorophore emission was demonstrated to change dramatically when it is placed on the different distances within PN, depending on the electric field. There are multiple factors that affect the molecule fluorescence signal: type of material, size and shape of the nanostructure, type of material and thickness of the nanoparticle coating, incident and emitted light wavelengths and intrinsic quantum yield of the fluorophore [27]. At the present time, scientists agree that there are two main factors affecting the fluorescence changes by plasmonic nanostructures. First of them is the plasmon field generated around PN by the incident light that, depending on wavelength, can enhance the excitation of the fluorophore, which, in turn, determines the level of fluorescence emission. The second one is the PN-fluorophore interaction that reduces the ratio of radiative to non-radiative decay rate and, depending on the presence of the dielectric layer, influences the quantum yield of the fluorophore, resulting in fluorescence quenching [28].

Several mathematical models have been developed to describe the plasmon electric field around PN [20, 29], including an approach considering dielectric-coated PN [27], which is of special importance for PEF modeling. It is based on the electrostatic approximation by solution of Laplace’s equation with the boundary conditions appropriate to the selected model [29] and takes into account only the dipole plasmon mode excited in spherical PN (Fig. 1).

The model shown in Fig. 1 uses the spherical coordinate system. According to this simplified approach [27], where the system has an azimuthal symmetry \( \frac{dE_P}{d\phi} = 0 \), plasmon electric field \( (E_P) \) at an observation point \( r \), which is generated by the incident light \( (E_0) \) in the vicinity of PN (with a radius \( r_1 \)) coated with a dielectric shell (with thickness \( r_2 - r_1 \)), can be described by the following equation

\[
\vec{E}_P = \left\{ \begin{array}{l}
2 \frac{\varepsilon_3 \varepsilon_a - \varepsilon_3 \varepsilon_b}{\varepsilon_2 \varepsilon_a + 2 \varepsilon_3 \varepsilon_b} \left( \frac{r_2}{r} \right)^3 + 1 \right\} \vec{E}_0 \cdot \cos(\theta) \cdot \hat{e}_r + \\
\left\{ \begin{array}{l}
\varepsilon_2 \varepsilon_a - \varepsilon_3 \varepsilon_b \\
\varepsilon_2 \varepsilon_a + 2 \varepsilon_3 \varepsilon_b
\end{array} \right\} \left( \frac{r_2}{r} \right)^3 + 1 \right\} \vec{E}_0 \cdot \sin(\theta) \cdot \hat{e}_\theta,
\end{array} \right.
\]

and the electric field inside the shell \( (\vec{E}_P^{layer}) \) is

\[
\vec{E}_P^{layer} = \left\{ \begin{array}{l}
3 \varepsilon_a \left( \varepsilon_1 + 2 \varepsilon_a \right) + 2 \left( \varepsilon_a - \varepsilon_1 \right) \left( \frac{r_2}{r} \right)^3 \right\} \\
\varepsilon_1 \varepsilon_a + 2 \varepsilon_2 \varepsilon_a
\end{array} \right\} \times \\
\times \vec{E}_0 \cdot \cos(\theta) \cdot \hat{e}_r - \left\{ \begin{array}{l}
3 \varepsilon_a \left( \varepsilon_1 + 2 \varepsilon_a \right) - 2 \left( \varepsilon_a - \varepsilon_1 \right) \left( \frac{r_2}{r} \right)^3 \\
\varepsilon_1 \varepsilon_a + 2 \varepsilon_2 \varepsilon_a
\end{array} \right\} \times \\
\times \vec{E}_0 \cdot \sin(\theta) \cdot \hat{e}_\theta.
\]

Fig. 1. Model of a dielectric-coated spherical PN placed in surrounding medium with a coordinate system used for the presentation of plasmon electric field.

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where \( \vec{e}_r, \vec{e}_\theta, \vec{e}_\phi \) are the unit vectors in \( r \) and \( \theta \) directions of the spherical coordinates, respectively,
\[
\varepsilon_\omega = \varepsilon_1(3 - 2P) + 2\varepsilon_2 P, \quad \varepsilon_b = \varepsilon_1 P + \varepsilon_2(3 - P),
\]
\[
P = 1 - \left( \frac{\varepsilon_2}{\varepsilon_1} \right)^2,
\]
\( \varepsilon_1, \varepsilon_2, \varepsilon_3 \) and \( \varepsilon_0 \) are the dielectric permittivity values of PN, shell, outer surrounding medium and vacuum, respectively. For PN, \( \varepsilon_1 \) is wavelength- and size-dependent and can be described by the Drude-Lorentz model [30]:
\[
\varepsilon_1(\omega) = \varepsilon_0 \left( 1 - \frac{\omega_p^2}{\omega^2 + i\omega\gamma_f} + \frac{\omega_p^2}{\omega_0^2 - \omega^2 + i\omega\gamma_b} \right),
\]
where \( \omega \) is the angular frequency of the incident light, \( \omega_0 \) – bound electron resonant frequency, \( \omega_p \) – plasma frequency,
\[
\gamma_f = \frac{1}{\tau_f} = \frac{1}{\tau_0} + v_f / \eta_f, \quad \gamma_b = \frac{1}{\tau_b},
\]
where \( \tau_f \) and \( \tau_b \) are the free electron relaxation time and bound electron decay time, respectively, \( v_f \) is the Fermi velocity and \( \tau_0 \) – free electron scattering time for the bulk material.

Quantum yield \( q \) indirectly influenced by the plasmon field \( E_p \) [27] can be described as
\[
\frac{q}{q_0} = \frac{\gamma_r}{\gamma_r^0} \left( \frac{\gamma_{abs}}{\gamma_r} + 1 - q^0 \right),
\]
where \( \gamma_r \) and \( \gamma_r^0 \) are the radiative decay rates in the presence and absence of PN, respectively, \( \gamma_{abs} \) is the non-radiative decay rate resulting from the energy absorbed by PN [31], and \( q^0 \) – intrinsic quantum yield of the fluorophore.

Therefore, fluorescence emission enhancement (\( \Phi \)) is the combined effect of the excitation rate enhancement and quantum yield change, both influenced by the plasmon field, and it can be described as
\[
\Phi = \left( \frac{E_p}{E_0} \right)^2 \frac{q}{q^0}.
\]

The electric field intensity near the 100-nm-diameter spherical gold nanoparticle with a 20-nm-thick SiO2 shell has been calculated using the above-mentioned theoretical approach implemented in the MathCAD software. The results of simulation are presented in Fig. 2 as the spectral distribution of average electric field intensity enhancement in the 3-nm vicinity around the spherical gold nanoparticle with a coating, where fluorophore molecules can reside. This system exhibits the enhancement of incident electric field intensity up to approximately 18 times with a maximum enhancement at the wavelength of 542 nm.

3. Sample preparation

Microscope glass slides (13×25 mm) were used as substrates. Before metal film deposition, substrates were cleaned in the ultrasonic bath with surfactant, then treated in “piranha solution” for 30 min (caution this mixture is potentially explosive and it will became very hot), triply rinsed with a copious amount of deionized water and finally dried using N2 flow.

Gold island film fabrication and SiO2 coating were carried out by thermal vacuum evaporation method using the UVN-2M setup (pressure 10^-3 Pa) with deposition speed of about 0.11 to 0.14 nm/s. The mass thickness of gold island film of about 10 nm was chosen as the value that allows to obtain the separated well-defined nanostructures after annealing (450 °C, 2 hours in air atmosphere). To prepare the dielectric coating, SiO2 layers with a thickness of 10 to 25 nm in 5-nm steps were deposited. Resulting structures (i.e. PN with the dielectric coating) were covered with fluorescent organic dye R6G by dipping them for 30 s into R6G aqueous solution with the dye concentration of 10^-3mol/L and subsequent drying under room conditions. The presence of R6G layer on glass or SiO2 surface was provided by opposite charges of SiO2 and R6G in aqueous solution [32, 34].

4. Experimental methods

Absorbance spectra of gold island films, PN without coating and PN with SiO2 coating were measured using LSPR spectrometer “NanoPLASMON-003” [34]. Atomic force microscopy (AFM) measurements of nanochip samples were carried out on atomic force microscope NanoScope IIIa Dimension 3000. Developed by the authors portable laser-based fluorometer

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“Fluorotest\textsuperscript{Nano}” was used for steady-state fluorescence measurements (see schematic measurement setup and common view of the device in Fig. 3). The device has three built-in laser light sources (405, 532, 650 nm) and four LEDs (360, 470, 535, 632 nm) that allows working both with nanochips and solutions in photometric cuvettes.

Fluorescence measurements of R6G layers on PN were carried out with the 532-nm green laser used as an excitation light source. The measurement scheme for fluorescence spectroscopy based on right-angle geometry with the 540-nm cut-off interference filter is depicted in Fig. 4. The simplified scheme of gold PN with the SiO\textsubscript{2} dielectric coating, which provides fluorescence signal enhancement for the adjacent dye molecule, is shown in Fig. 4. The distance between the fluorophore molecules and PN was determined by the thickness of the dielectric SiO\textsubscript{2} coating deposited on PN.

5. Results and discussion

The top view AFM image of the gold island film produced according to the protocol described in Section 3 is shown in Fig. 5a. AFM measurements demonstrate that the maximal mass thickness of island film is about 12 nm, which is close to the value specified in the fabrication protocol. Gold PN (Fig. 5b) after the thermal annealing of gold island film exhibit the shape of semi-ellipsoids with a height and an equivalent radius distributed within the ranges of 30…80 nm and 20…100 nm, respectively. The average height of PN is about 60 nm and the average equivalent radius is close to 50 nm.

Typical absorbance spectra of bare gold PN and gold PN with SiO\textsubscript{2} coating of a different thickness (10…25 nm) measured in air are shown in Fig. 6. All spectra reveal expressed plasmonic peak at 550 nm (for bare PN) and 580…610 nm (for covered PN). The observed LSPR position red shift accompanying the dielectric coating thickness increase confirms the plasmonic nature of registered absorbance peak [35].
Fig. 5. Top view AFM images of (a) gold island film and (b) gold PN.

Fig. 6. Absorbance spectra of bare gold PN and gold PN with SiO₂ dielectric coating having the thickness within the range of 10...25 nm.

All the studied PN samples covered with SiO₂/dye exhibit R6G fluorescence enhancement in comparison with the signal obtained from similar samples without PN (Fig. 7a). At the same time, the relationship between the intensity of fluorescent signal and SiO₂ thickness is evidently nonlinear with a distinct peak. Namely, maximum R6G fluorescence was obtained for the sample with a 20-nm-thick SiO₂ coating. Fig. 7b shows the fluorescence enhancement factor dependence on the thickness of dielectric coating, which exhibits the bell-shaped character. Here, enhancement factor is estimated as a ratio of R6G fluorescence intensities for gold PN with SiO₂ shell and bare glass substrate near maximum emission wavelength (548 nm). Small blue-shift of R6G emission with respect to tabular value can be explained by presence of silicon oxide shell near the dye, which changes the dielectric environment of R6G [36].

Fig. 7. (a) Fluorescence spectra of R6G dye placed on the surface of gold PN coated with SiO₂ layer of various thicknesses. (b) R6G fluorescence enhancement factor dependence on the SiO₂ thickness.
It is of importance that the enhanced signal was observed only if the proximity of wavelengths of PN LSPR, dye emission and dye excitation has been provided, which is in agreement with well-known phenomenon of emission enhancement by strong electric field generated by PN under LSPR conditions [37]. The observed increase of dye emission using the proposed nanochips opens the possibility to provide plasmon-enhanced fluorescence for flow mode real-time measurements with low analyte consumption. The experimental results of R6G fluorescence enhancement close to 22 times, obtained on the fabricated nanochips with 20-nm SiO$_2$ coating, are in agreement with the results of modeling. It should be noted that the measured and calculated absorbance spectra of PN have some difference in the wavelength positions. This disagreement can be attributed to the presence of the dielectric substrate and the different sizes and shape of PN on the real nanochip. In particular, nanoparticles on the nanochip have rather semi-ellipsoid, not a spherical shape. Disagreement in enhancement values obtained in the model and the experimental data can also be caused by a simplified model “nanoparticle – dielectric coating – dye” In the applied model, the quantum yield of the dye that affects the value of the energy transfer between the LSPR nanoparticle and dye molecules has not been considered, as well as interaction between nanostructures in array, this will be the part of our future work.

6. Conclusions

The method of fluorescence signal enhancement based on exploitation of plasmon-supporting nanochips was developed, which provides up to 22-fold emission enhancement for Rhodamine 6G. The relevant experimental studies have been carried out on the novel laser-based fluorometer “FluorotestNano” that provides plasmonic control of fluorescence signal in the visible spectral region. Nanochips for plasmon-enhanced fluorescence based on gold nanostructures with dielectric coating of different thickness were designed and fabricated by thermal annealing of gold island films with subsequent SiO$_2$ layer deposition. As a result of R6G fluorescence intensity measurements for different thicknesses of dielectric spacer, the optimal distance between dye molecules and plasmonic nanostructures providing a dominance of fluorescence enhancement over quenching was found. The obtained results demonstrate that the creation of “gold nanostructure – dielectric spacer – fluorophore” system with precisely tailored optical and geometrical parameters and distinct resonant properties enables the development of high-sensitive optoelectronic nanosensors for the analysis of weak fluorescence signals, e.g. for biochemical and diagnostic applications.

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References


