

## Bilayer MoS<sub>2</sub> phototransistor deposited on a flexible substrate

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**Abstract.** A new phototransistor with a MoS<sub>2</sub> bilayer deposited on a polyimide substrate was fabricated using pulsed-laser deposition technique. The structural and optical properties of the bilayer MoS<sub>2</sub> films were characterized by Raman, photoluminescence and optical absorption spectroscopy. The photocurrent generated by a laser light with the wavelength of 532 nm in the top gate flexible field effect transistor was determined for specified drain and gate voltages. The photoresponsivity of the device could be adjusted over a wide range by controlling the gate and drain voltages. At this, the photoresponsivity value could exceed 12.5 mA/W. The structure showed switching behavior with a characteristic time of approximately 10 ms in response to pulsed illumination. The photocurrent changed significantly due to compressive or tensile strain induced by the bent substrate, which affected the band-gap value of the ultrathin MoS<sub>2</sub> layers. The proposed structure may find potential applications in flexible electronics, such as *e.g.* for nanoscale flexible photodetectors or stress sensors.

**Keywords:** bilayer MoS<sub>2</sub>, flexible substrate, phototransistor, strain.

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### 1. Introduction

Recently, a great interest has been drawn to 2D semi-conducting materials such as MoS<sub>2</sub>. This transition-metal dichalcogenide exhibits interesting electrical, optical and mechanical properties that correlate with its layered structure [1].

Using band and strain engineering, various functional 2D heterostructures for applications in high-performance optoelectronics, such as *e.g.* for photodetectors covering a wide spectral range from infrared to ultraviolet one [2–6] can be designed.

It is well known that a layered MoS<sub>2</sub> crystal consists of triple-layered sheets with hexagonally distributed Mo atoms in the middle, which have strong covalent bonding to six sulfur atoms in the top and bottom layers. At this, the inter-sheet stacking is maintained by the van der Waals forces [1]. Such weak interaction between the neighboring sheets having no surface dangling bonds enables growing MoS<sub>2</sub> on various substrates without need to account for possible lattice mismatch. Such substrates may include organic polymers, in particular polyimide (PI). Direct growth of 2D materials on flexible substrates is a big step toward fabrication of 2D flexible devices [7–9].

Numerous fabrication approaches, including sputtering, chemical vapor deposition and molecular beam epitaxy, have been used to grow 2D materials.

Unfortunately, most of the synthesis processes, in particular, the widely exploited chemical vapor deposition, use the growth temperatures above 500 °C and therefore are not compatible with polymers that have low thermal stability [10, 11].

In contrast, the temperature during pulsed laser deposition (PLD) is relatively low, enabling direct growth of 2D materials on different plastic substrates. The PLD also offers other advantages for preparing various layered 2D materials and structures on a large scale [12–14]. It can be used to transfer the target material to different substrates according to their stoichiometric ratio. In this method, the layer imperfection and the band gap of the 2D material can be efficiently controlled by simply manipulating the energy and number of laser pulses.

It is well established that optical properties of 2D layered semiconductors are highly controlled by the number of the layers owing to the quantum confinement effects in the out-of-plane direction and changes in symmetry [15, 16]. Despite their high (more than 97% [17]) optical transparency, mono- or few-layer 2D materials can strongly interact with incident light, which leads to enhanced photon absorption. This absorption stems from existence of the Van Hove singularities in the electronic density of states and formation of strongly-bound excitons/trions induced by the reduced thicknesses [18–20].

Moreover, atomic thickness of 2D layered materials provides extraordinary mechanical flexibility. Flexible bending of substrates can cause mechanical strain in ultrathin layers, which changes the inter-atomic distances. Such changes can impact material properties such as band gap value, optical absorption coefficient and carrier mobility [4, 9, 19, 21]. Note that the strain effects may be significant for 2D materials without any bending. For example, the thermal expansion coefficient of 2D MoS<sub>2</sub> highly exceeds that of plastic or rigid substrates [22], which causes a tension in the thin film during post-growth cooling.

It has been demonstrated that atomically thin MoS<sub>2</sub> is a promising channel material for field-effect phototransistors (FET) with a high current ON/OFF ratio, a low sub-threshold slope, and high photoresponsivity in the wide spectral range [23].

Hence, direct synthesis of high-quality MoS<sub>2</sub> on plastic substrates is of high demand for flexible and wearable optoelectronic devices.

The main objectives of this study are to fabricate phototransistors based on bilayer (2L) MoS<sub>2</sub> on PI substrates using the PLD method and to make a detailed analysis of their dark and light-induced electric properties.

## 2. Experimental results and discussions

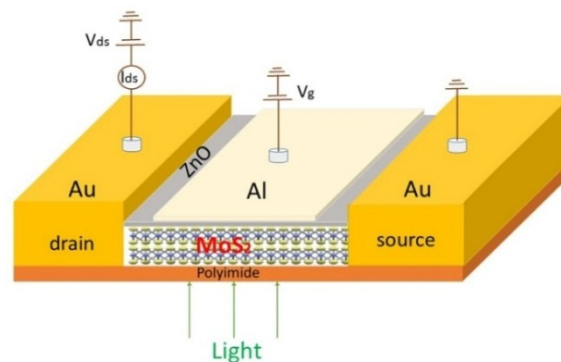
### 2.1. Preparation of 2L MoS<sub>2</sub> phototransistor

A Q-switched Nd laser was used to deposit thin films consisting of a few atomic monolayers of MoS<sub>2</sub>. The laser operated at a wavelength of 1.064  $\mu\text{m}$ , pulse duration of 30 ns, pulse energy of 0.35 J, pulse repetition rate of 0.1 Hz, and beam diameter of 20 mm. A 50  $\mu\text{m}$  thick polyimide film was chosen as a flexible substrate due to its exceptional physical, electrical and mechanical properties, which remain stable over a wide temperature range (up to 400 °C).

Ultrathin MoS<sub>2</sub> layers were grown by DC magnetron sputtering using a target in form of a pressed tablet of MoS<sub>2</sub> and sulfur powders taken in such weight ratio that the atomic ratio of Mo:S in the target was 1:4 [13]. For the growth process, the target and the substrate were located in the growth chamber at a distance of 3 to 6 cm from each other. The vacuum level in the chamber was maintained no worse than  $10^{-5}$  mm Hg. The substrate was heated to 200 °C.

After the growth, the films were annealed at 200 °C for 5 min and then cooled down to room temperature. The thickness of the MoS<sub>2</sub> films can be controlled by the number of pulses used to grow the films. Typically, 5 laser pulses are sufficient to grow one monolayer of MoS<sub>2</sub> under the used technological conditions [13]. The grown films had good adhesion to the polyimide substrate.

Prior to fabricating a MoS<sub>2</sub> phototransistor, we characterized the deposited 2L MoS<sub>2</sub> films by measuring their photoluminescence (PL), optical transmittance and Raman spectra.



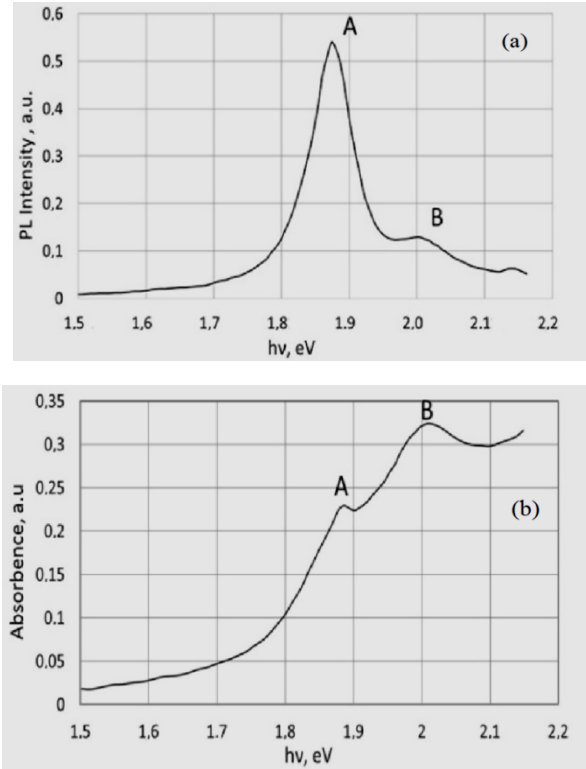
**Fig. 1.** Schematic view of a 2L MoS<sub>2</sub> FET structure with electrical contacts. The direction of incident light used to characterize the device is indicated.

We fabricated a FET using 2L MoS<sub>2</sub> as the conductive channel and a 100 nm thick transparent *i*-ZnO (high- $k$  dielectric) layer as the top-gate dielectric. This allows for low power consumption due to an increased gate capacitance and enhanced mobility caused by the dielectric screening effect [24]. 100 nm thick gold layers deposited on the PI substrate prior to the 2L MoS<sub>2</sub> film deposition were used as the source and drain. An Al gate electrode was deposited onto the gate dielectric (see Fig. 1). The gate length and width were 5 nm and 1  $\mu\text{m}$ , respectively. The electrical transport properties of the 2L MoS<sub>2</sub> FET were measured in the dark and under illumination of a green laser with  $\lambda = 532$  nm at powers of up to 70  $\mu\text{W}$  at room temperature.

### 2.2. Structure characterization

During a standard experiment, an ultrathin MoS<sub>2</sub> layer was directly deposited onto a polyimide substrate along with gold drain and source electrodes. The film thickness controlled by the number of ablating laser pulses was approximately 1.2 nm. The film consisted of two monolayers as was identified by Raman spectroscopy investigations.

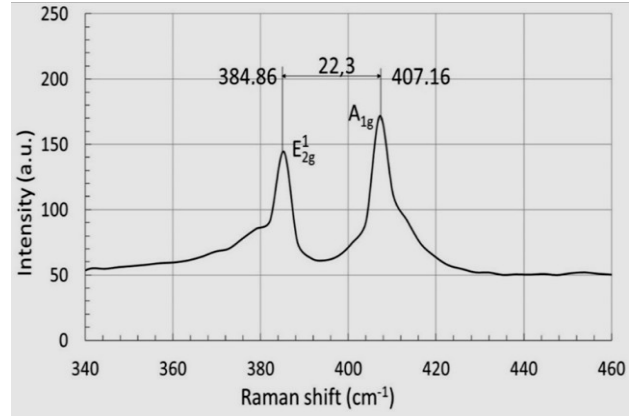
The PL spectra of the deposited films were measured at room temperature using an excitation laser with a wavelength of 532 nm. It can be seen from Fig. 2a that the PL spectrum shows two clear emission peaks A and B at 1.89 and 2.01 eV, respectively. These peaks are related to direct radiative recombination of photo-generated electron-hole pairs that form two types of excitons corresponding to the spin-orbital split valleys of the valence band at the K-point of the Brillouin zone [25, 26]. The peak energy separation (approximately 0.11 eV) corresponds to spin-orbital splitting of the valence band at the K-point of the Brillouin zone. An apparently dominant peak A points to a low density of non-radiative recombination defects and high sample quality [26]. The asymmetric shape of the emission peak A may be caused by trion contribution, which can be more prominent in 2L MoS<sub>2</sub> as compared to 1L MoS<sub>2</sub> [20, 22].



**Fig. 2.** PL (a) and optical absorbance (b) spectrum of 2L MoS<sub>2</sub> grown on PI substrate by PLD (A – 1.88 eV, B – 2.02 eV).

Optical absorbance of the 2L MoS<sub>2</sub> as a function of photon energy is presented in Fig. 2b. It can be seen from this figure that the film is characterized by high absorbance in the visible range. In particular, absorbance for the photon energies greater than 1.8 eV exceeds 10%. The absorption peaks A and B at 1.89 and 2.01 eV in Fig. 2b match the PL resonances in both position and width. Therefore, they may be attributed to direct excitonic transitions between the spin-orbital split valence band maxima and conduction band minimum, all located at the K-point of the Brillouin zone. The absorption (emission) tail in the photon energy range of 1.5 to 1.8 eV is due to indirect transitions between the valence band maximum and the conduction band minimum at the  $\Gamma$ -point of the Brillouin zone [15, 25, 26].

Since Raman spectroscopy is a powerful tool for determining vibrational modes of atomically thin 2D materials, we used it to further describe the deposited MoS<sub>2</sub> films and verify their thickness and properties (Fig. 3). Two typical Raman active modes of hexagonal MoS<sub>2</sub> at 385 and 407 cm<sup>-1</sup> were observed, which can be attributed to in-plane  $E_{2g}^1$  and out-of-plane  $A_{1g}$  lattice vibrations, respectively [27]. Separation between the  $E_{2g}^1$  and  $A_{1g}$  peaks depends on the number of MoS<sub>2</sub> single layers, which makes it possible to determine the number of the layer in the synthesized samples from the Raman spectra [27, 28]. Using an empirical relation for the variation of this interpeak separation with the number of the monolayers in MoS<sub>2</sub> thin films [13, 27], this value is found to be about 22.6 cm<sup>-1</sup> for the bilayer MoS<sub>2</sub>.



**Fig. 3.** Raman spectrum of 1.2 nm thick MoS<sub>2</sub> film grown on PI substrate by PLD.

This value is very close to the measured value of 22.3 cm<sup>-1</sup> shown in Fig. 3. The narrowness of the main Raman peak  $A_{1g}$  (FWHM = 4 cm<sup>-1</sup>) is indicative of the high crystalline quality of the synthesized film. The small difference between these values may be attributed to the possible effect of the strain caused by the mismatch between the thermal expansion coefficients of the PI substrate and deposited MoS<sub>2</sub> layer [21, 29].

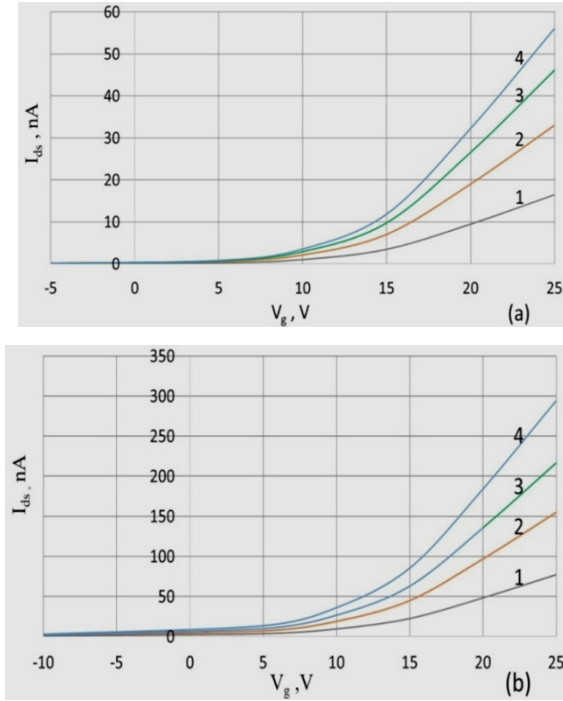
The MoS<sub>2</sub>-based phototransistor was investigated by studying its output characteristics, such as generated photocurrent and photoresponsivity, as well as its switching behavior at different values of optical power, and gate and drain voltage.

The effect of the mechanical strain exerted on the bent flexible substrate was also studied. The gating (transfer) characteristics of our FET are shown in Fig. 4 confirming that it is a typical FET with an *n*-type channel. Varying the gate voltage  $V_g$  from negative to low positive values does not significantly affect the characteristics, thus indicating that the top gate is unlikely to accumulate trapped charges during the experiment.

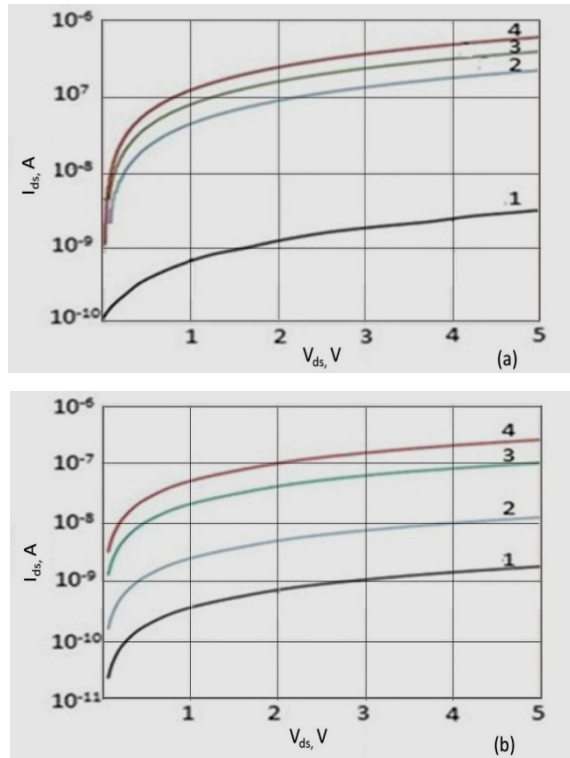
### 2.3. Photocurrent (PC)

As the device was illuminated by a green laser with the wavelength  $\lambda = 532$  nm at room temperature, the drain current increased by several orders of magnitude. The output curves in the dark and under illumination are shown in Fig. 5. It can be seen from this figure that our transistor exhibits current saturation at high drain voltages  $V_{ds}$ . The saturation current under the “pinch-off” conditions increases with the illumination power and also depends on the back gate voltage.

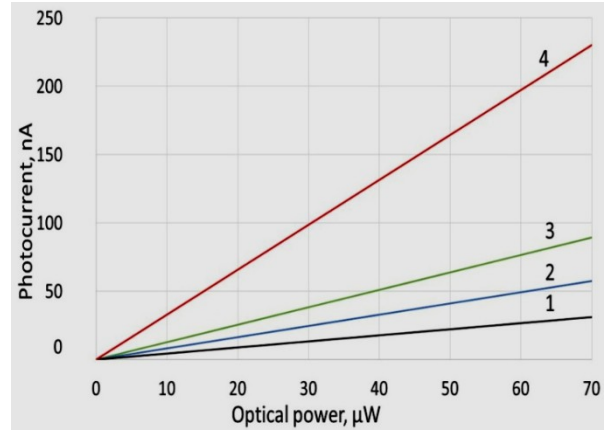
The device photoresponsivity ( $R_{ph}$ ) defined as  $I_{ph}/P_{light}$  (here,  $I_{ph}$  is the difference between the PC and the dark current and  $P_{light}$  is the total optical power, respectively) depends on the drain and gate voltages. For example, at  $V_g = 0$  V and  $V_{ds} = 4.5$  V,  $R_{ph}$  is approximately equal to 5 mA/W (Fig. 6). When the gate voltage increases to 40 V,  $R_{ph}$  becomes higher than 12.5 mA/W, thus proving that the top gate plays an important role in controlling photocurrent in the 2L MoS<sub>2</sub> phototransistor.



**Fig. 4.** Gating response (transfer characteristics) of MoS<sub>2</sub> phototransistor in the dark (a) and under illumination at a power of 10 μW (b), acquired at the top gate voltage  $V_g$  varying from -10 to +25 V and at different drain voltages  $V_{ds} = 1$  V (1), 2 (2), 3 (3), and 4 (4).



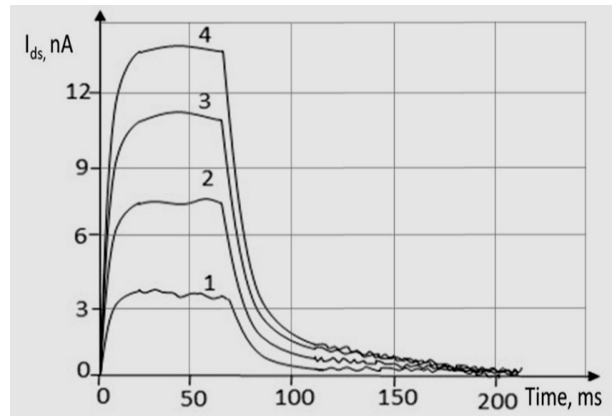
**Fig. 5.** (a) Output ( $I_{ds}$ - $V_{ds}$ ) characteristics of the device at  $V_g = 0$  V in the dark and under illumination at different power values: in the dark (1), 0.5 (2), 1 (3), and 5 μW (4). (b) ( $I_{ds}$ - $V_{ds}$ ) characteristics under illumination (5 μW) at different gate voltages  $V_g = 0$  (1), 10 (2), 20 (3), and 30 V (4).



**Fig. 6.** PC of FET as a function of illumination power measured at  $V_g = 0$  V and different drain voltages  $V_{ds} = 0.5$  (1), 1 (2), 1.5 (3), and 4 V (4). The device is illuminated from the PI substrate side.

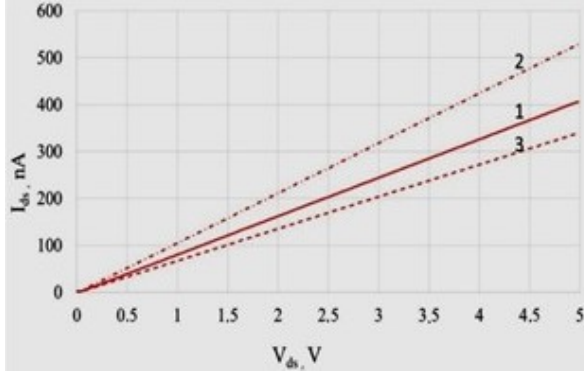
As shown in Fig. 6, the PC measured at  $V_g = 0$  V and at different drain voltages is proportional to the illumination power due to an increase in the electron-hole pair generation rate caused by light absorption in the MoS<sub>2</sub> layer.

A sharp rise/decay of the PC was observed when the radiation was either turned ON or OFF for all the drain and gate voltages (Fig. 7). The response time (defined as the time to reach 90% of the maximum PC after switching the light from OFF to ON) and the recovery time (defined as the time to reach 10% of the maximum PC after switching light from ON to OFF) were estimated to be 12 and 25 ms, respectively. Such an abrupt rise and decay of the PC observed for different values of the drain (gate) voltage and illuminating light power can be explained by the lack of metastable charge carrier trapping centers in the bulk of the MoS<sub>2</sub> and at the interfaces.



**Fig. 7.** Dynamical photo-switching behavior of phototransistor at  $V_g = 0$  V and  $V_{ds} = 4$  V for different illumination power  $P_{light}$  (μW): 10 (1), 20 (2), 30 (3), and 40 (4). The laser is turned on for 75 ms and then turned off.





**Fig. 8.** PC as a function of  $V_{ds}$  before the substrate bending (1) and after its convexly (2) or concavely (3) curving ( $\varepsilon = \pm 0.5\%$ ). The curving of the substrate induces a tensile (2) or compressive stress (3) in the film, respectively. The gate voltage  $V_g = 0$  V and the illumination power is  $1 \mu\text{W}$ .

To evaluate the influence of bending on the MoS<sub>2</sub>-based phototransistor, PC as a function of the drain voltage was measured before and after the bending. In general, as shown in Fig. 8a, the PC can significantly change even when a slight stress is applied. In the experiment, the flexible MoS<sub>2</sub> FET was placed in bending molds with convex or concave curvatures. The uniaxial strain applied to the sample, estimated as the ratio of the substrate thickness ( $h_{PI}$ ) to twice the radius ( $R$ ) of curvature ( $\varepsilon \approx h_{PI}/2R$ ), was  $0.5\%$  for  $h_{PI} = 50 \mu\text{m}$  and  $R = 5$  mm. The mechanical stress  $\sigma_{bend}$  in the MoS<sub>2</sub> film related to the bending strain can be estimated using the following expression [30]:

$$\sigma_{bend} = \frac{E_{PI} h_{PI}^2}{6d_{\text{MoS}_2} R (1 - \nu_{PI})}, \quad (1)$$

where  $E_{PI}$  and  $\nu_{PI}$  are the Young modulus and the Poisson ratio of the substrate, respectively, and  $d_{\text{MoS}_2}$  is the MoS<sub>2</sub> film thickness. The value of  $E_{PI}$  used in the calculations was  $2.5$  GPa for polyimide films (as indicated by the manufacturer [31]) and  $\nu_{PI}$  was fixed at  $0.34$ . The estimated value of  $\sigma_{bend}$  was  $263$  GPa.

In general, the total stress in a 2L MoS<sub>2</sub> is composed of the stress induced by bending and the thermal stress originating from a mismatch in the coefficients of thermal expansion (CTE) between the PI substrate and the thin MoS<sub>2</sub> film. Polyimide films are characterized by the CTE value  $\alpha_{PI} = 11 \cdot 10^{-6} \text{ 1/C}$  [31], which is usually much larger than that of 2L MoS<sub>2</sub> ( $\alpha_{\text{MoS}_2} = 0.5 \cdot 10^{-6} \text{ 1/C}$  [29]). Therefore, during cooling down from the elevated temperature corresponding to the deposition process (approximately  $200^\circ\text{C}$ ) to room temperature, the PI film shrinks at a high rate, which results in a biaxial compressive stress in the MoS<sub>2</sub>. In other words, the MoS<sub>2</sub> films grown on the PI substrates were always under a certain compressive stress even without any bending. The thermal stress ( $\sigma_{th}$ ) can be expressed as follows [32]:

$$\sigma_{th} = E_{\text{MoS}_2} (E_{PI} - E_{\text{MoS}_2}) (T_{dep} - T_{measure}), \quad (2)$$

where  $E_{\text{MoS}_2}$ ,  $\alpha_{\text{MoS}_2}$ ,  $\alpha_{PI}$ ,  $T_{dep}$ , and  $T_{measure}$  are the Young modulus of MoS<sub>2</sub>, CTEs of MoS<sub>2</sub> and the polyimide film, and the sample temperatures during the deposition and the measurement, respectively. It is noteworthy that the Poisson effect is neglected in Eq. (2). The temperatures  $T_{dep}$  and  $T_{measure}$  were set to  $200^\circ\text{C}$  and  $25^\circ\text{C}$ , respectively, and  $E_{\text{MoS}_2}$  for 2L MoS<sub>2</sub> was set to  $270$  GPa [21, 33]. The value of  $\sigma_{th}$  estimated by Eq. (2) is  $0.74$  GPa, which is a negligible value compared to the bending induced stress.

We may conclude therefore that when the PI film is convexly curved upward, the MoS<sub>2</sub> film deposited on its upper surface is subjected to a tensile stress. Conversely, when the MoS<sub>2</sub> film is deposited on the concave side of a bent PI substrate, it is subjected to a compressive stress. Convex bending, as is known, results in a reduction of the band-gap of 2L MoS<sub>2</sub> due to the tensile strain [21], while concave bending leads to an increase in the band gap of the MoS<sub>2</sub> [34, 35]. A decrease in the band gap led to the increased light absorption by the MoS<sub>2</sub>. As a result, for a given laser intensity, and the drain and gate voltage, the structure exhibited a larger photoresponse compared to the case of a flat substrate. Under the same conditions, concave bending (compressive strain), increasing the band gap, should lead to a decreased light absorption and, therefore, lower device photocurrent as shown in Fig. 8.

### 3. Conclusions

In summary, employing the low-temperature PLD method, we fabricated a 2L MoS<sub>2</sub>-based flexible phototransistor and conducted a detailed investigation of its electric and photoelectric characteristics. Under fixed drain and gate voltages, the PC is depended only on the illumination intensity. The photoresponsivity of the device can be adjusted over a wide range by controlling the gate and drain voltages and can be higher than  $12.5 \text{ mA/W}$ . This parameter for the flexible 2L MoS<sub>2</sub> phototransistor fabricated on a polyimide substrate is not inferior to that of a phototransistor based on graphene or 1L MoS<sub>2</sub> [2, 3, 5].

We have shown that generation and recombination of non-equilibrium carriers in our device can be switched within tens of ms. Such prompt photo-switching behavior controlled by incident light highlights the stability of the device characteristics.

We observed a pronounced change in the PC due to bending-induced strain variation in the band gap. This observation confirms the viability of strain engineering as a powerful tool for creating new flexible and tunable photodetectors based on several-layer MoS<sub>2</sub> deposited on PI substrates. Among the potential applications of such few-layer MoS<sub>2</sub> structures there are nanoscale stress sensors and displays.

### Declaration of competing interest

The authors declare no competing financial interests or personal relationships to the reported paper.

### Data availability

The reported data can be available from the corresponding authors upon request.

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#### Фототранзистор на основі двошарового MoS<sub>2</sub>, нанесений на гнучку підкладку

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**Анотація.** Новий фототранзистор на основі двошарового MoS<sub>2</sub>, нанесеного на поліімідну підкладку, виготовлено методом імпульсного лазерного осадження. Структурні та оптичні властивості двошарових плівок MoS<sub>2</sub> охарактеризовано, використовуючи раманівську спектроскопію, фотолюмінесценцію та спектроскопію оптичного поглинання. Фотострум, що генерується лазерним світлом з довжиною хвилі 532 нм у гнучкому польовому транзисторі з верхнім затвором, визначено для заданих напруг стоку та затвора. Фоточутливість пристрою можна регулювати в широкому діапазоні, задаючи значення напруги затвора та стоку, і вона може перевищувати 12,5 мА/Вт. Під дією імпульсного освітлення структура проявляє властивість перемикачання з характерним часом близько 10 мс. Фотострум значно змінюється через деформацію стискання або розтягу, викликану вигнутістю підкладки, що впливає на ширину забороненої зони в ультратонких шарах MoS<sub>2</sub>. Запропонована структура може знайти потенційне застосування для створення гнучких електронних приладів, таких як нанорозмірний гнучкий фотодетектор або датчик напруги.

**Ключові слова:** двошаровий MoS<sub>2</sub>, гнучка підкладка, фототранзистор, деформація.